

"But our present QM formalism is not purely epistemological; it is a peculiar mixture describing in part realities of Nature, in part incomplete human information about Nature --- all scrambled up by Heisenberg and Bohr into an omelette that nobody has seen how to unscramble. Yet we think that the unscrambling is a prerequisite for any further advance in basic physical theory. For, if we cannot separate the subjective and objective aspects of the formalism, we cannot know what we are talking about; it is just that simple." --E.T. Jaynes

Formulating Quantum Theory as a Causally Neutral Theory of Bayesian Inference

Robert Spekkens Perimeter Institute

Joint work with Matt Leifer Iniversity College London

APCWQIS, Putrajaya, December 2012

Where I'm coming from...



In the resulting model quantum states are states of incomplete knowledge



Epistemically Restricted Liouville mechanics

A Liouville distribution $\mu(q,p)$, can describe an agent's knowledge only if it satisfies:

The classical uncertainty principle:

 $\Delta^2 q \Delta^2 p - C_{q,p}^2 \geq (\hbar/2)^2$

The max-ent condition: the entropy of $\mu(q,p)$, $S(\mu) = -\int \mu(q,p) \log \mu(q,p) dq dp$

is maximal for a given set of second-order moments.



(This can be generalized to *n* systems)

The limit of perfect knowledge for some variables

An observer can only jointly know the values of a set of variables if they commute relative to the Poisson bracket.



know $Q_A - Q_B$ and $P_A + P_B$

$$P(q,p) \propto \delta(q-a) \qquad P(q,p) \propto \delta(p-b)$$

$$(q,p) \propto \phi(p-b)$$

$$P(q_A, p_A, q_B, p_B)$$

\$\propto \delta(q_A - q_B)\delta(p_A + p_B)\$

corresponds to EPR state

Valid transformations



Valid measurements



Theorem: Epistemically restricted Liouville mechanics is empirically equivalent to Gaussian quantum mechanics

Bartlett, Rudolph and Spekkens, arXiv:1111:5057

acausal connection



















Categorizing quantum phenomena

Those arising in a restricted statistical classical theory	Those not arising in a restricted statistical classical theory
Noncommutativity Entanglement Collapse Wave-particle duality Teleportation No cloning Interference Key distribution Improvements in metrology Quantum eraser Coherent superposition Pre and post-selection effects Others	Bell inequality violations Computational speed-up (if it exists) Noncontextuality inequality violations Certain aspects of items on the left Others

Quantum theory as a theory of Bayesian inference

Classical

P(R) Probability distributionover phase-spacecoordinates R

$$P(R = r)$$
 probability that *R=r*

Quantum

 ρ_A Operator on Hilbert space of A

No analogue (yet)

	Classical	Quantum
State of knowledge	P(R)	$ ho_A$
Normalization	$\sum_{R} P(R) = 1$	$\mathrm{Tr}_A \rho_A = 1$
Joint state	P(R,S)	$ ho_{AB}$
Marginalization	$P(S) = \sum_{R} P(R, S)$	$\rho_B = \mathrm{Tr}_A \rho_{AB}$
Conditional state	P(S R)	????

P(S|R)

Normalization condition

 $\sum_{S} P(S|R) = 1$

Conditional state

 $\rho_{B|A}$

Normalization condition

 $\operatorname{Tr}_B(\rho_{B|A}) = I_A$

P(S|R)

Normalization condition

 $\sum_{S} P(S|R) = 1$

Conditional state

 $ho_{B|A}$

Normalization condition $\operatorname{Tr}_B(\rho_{B|A}) = I_A$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$

Relation of conditional to joint

$$\rho_{B|A} = (\rho_A^{-1/2} \otimes I_B) \rho_{AB} (\rho_A^{-1/2} \otimes I_B)$$

P(S|R)

Normalization condition

 $\sum_{S} P(S|R) = 1$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$

Conditional state

 $ho_{B|A}$

Normalization condition $\operatorname{Tr}_B(\rho_{B|A}) = I_A$

Relation of conditional to joint $\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$

See: Leifer, PRA 74, 042310 (2006)

P(S|R)

Normalization condition

 $\sum_{S} P(S|R) = 1$

Relation of conditional to joint

 $P(S|R) = \frac{P(R,S)}{P(R)}$ P(R,S) = P(S|R)P(R)

Conditional state

 $\rho_{B|A}$

Normalization condition $\operatorname{Tr}_B(\rho_{B|A}) = I_A$

Relation of conditional to joint

$$\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$$
$$\rho_{AB} = \rho_A^{1/2} \rho_{B|A} \rho_A^{1/2}$$

P(S|R)

Normalization condition

 $\sum_{S} P(S|R) = 1$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$
$$P(R,S) = P(S|R)P(R)$$

Classical belief propagation

 $P(S) = \sum_{R} P(S|R)P(R)$

Conditional state

 $ho_{B|A}$

Normalization condition $Tr_B(\rho_{B|A}) = I_A$

Relation of conditional to joint

$$\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$$
$$\rho_{AB} = \rho_A^{1/2} \rho_{B|A} \rho_A^{1/2}$$

Quantum belief propagation $\rho_B = \text{Tr}_A(\rho_{B|A}\rho_A)$

Two formulas for the joint probability P(R,S) = P(S|R)P(R)= P(R|S)P(S)

Classical Bayes' theorem

 $P(S|R) = \frac{P(R|S)P(S)}{P(R)}$

Two formulas for the joint state $\rho_{BA} = \rho_A^{1/2} \rho_{B|A} \rho_A^{1/2}$ $= \rho_B^{1/2} \rho_{A|B} \rho_B^{1/2}$

Quantum Bayes' theorem

 $\rho_{B|A} = \rho_A^{-1/2} \rho_B^{1/2} \rho_{A|B} \rho_B^{1/2} \rho_A^{-1/2}$

Classical Bayesian conditioning

 $P(S) \to P(S|X = x)$ $P(S|X = x) = \sum_{X} P(S|X) \delta_{X,x}$

Quantum Bayesian conditioning

 $\rho_B \to \rho_{B|X=x}$ $\rho_{B|X=x} \equiv \operatorname{Tr}_X(\rho_{B|X}|x\rangle\langle x|_X)$

Causal Neutrality

$$\mathbf{R}$$
 \mathbf{S} $P(R,S)$ \mathbf{A} \mathbf{P}_{AB} $P(S|R) = P(R,S)/P(R)$ $????$ $P(S) = \sum_{R} P(S|R)P(R)$ $????$ $P(S) = \Gamma_{R \to S}[P(R)]$ $????$

S

R

$$P(R,S) \qquad ????$$

$$P(S|R) = P(R,S)/P(R) \qquad \textbf{B} \qquad ????$$

$$P(S) = \sum_{R} P(S|R)P(R) \qquad \textbf{A} \qquad ????$$

$$P(S) = \Gamma_{R \to S}[P(R)] \qquad \rho_{B} = \mathcal{E}_{A \to B}(\rho_{A})$$

$$\mathbf{R}$$
 \mathbf{S} $P(R, S)$ \mathbf{A} \mathbf{P}_{AB} $P(S|R) = P(R, S)/P(R)$ $\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$ $P(S) = \sum_R P(S|R)P(R)$ $\rho_B = \operatorname{Tr}_A(\rho_{B|A}\rho_A)$ $P(S) = \Gamma_{R \to S}[P(R)]$ $\rho_B = \mathfrak{E}_{A \to B}(\rho_A)$

$$P(R, S) \qquad ????$$

$$P(S|R) = P(R, S)/P(R) \qquad \blacksquare \qquad ????$$

$$P(S) = \sum_{R} P(S|R)P(R) \qquad \blacksquare \qquad ????$$

$$P(S) = \Gamma_{R \to S}[P(R)] \qquad \rho_{B} = \mathcal{E}_{A \to B}(\rho_{A})$$

R

$$\mathbf{R}$$
 \mathbf{S} $P(R, S)$ \mathbf{A} \mathbf{P}_{AB} $P(S|R) = P(R, S)/P(R)$ $\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$ $P(S) = \sum_R P(S|R)P(R)$ $\rho_B = \operatorname{Tr}_A(\rho_{B|A}\rho_A)$ $P(S) = \Gamma_{R \to S}[P(R)]$ $\rho_B = \mathfrak{E}_{A \to B}(\rho_A)$

$$\begin{split} P(R,S) & \varrho_{AB} \\ P(S|R) &= P(R,S)/P(R) & B & \varrho_{B|A} = \rho_A^{-1/2} \varrho_{AB} \rho_A^{-1/2} \\ P(S) &= \sum_R P(S|R)P(R) & A & \rho_B = \operatorname{Tr}_A(\varrho_{B|A} \rho_A) \\ P(S) &= \Gamma_{R \to S}[P(R)] & \rho_B = \mathcal{E}_{A \to B}(\rho_A) \end{split}$$

S

R

$$\mathbf{R}$$
 \mathbf{S} $P(R,S)$ \mathbf{A} \mathbf{P}_{AB} $P(S|R) = P(R,S)/P(R)$ $\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$ $P(S) = \sum_R P(S|R)P(R)$ $\rho_B = \operatorname{Tr}_A(\rho_{B|A}\rho_A)$ $P(S) = \Gamma_{R \to S}[P(R)]$ $\rho_B = \mathfrak{E}_{A \to B}(\rho_A)$ $\rho_{B|A} \ge 0$ $\mathfrak{E}_{A \to B} \circ T_A$ is CP

 α

$$\mathbf{R} - \cdots - \mathbf{S} \quad P(R, S)$$

$$P(S|R) = P(R, S)/P(R)$$

$$P(S) = \sum_{R} P(S|R)P(R)$$

$$\mathbf{A} - \cdots - \mathbf{B} \quad \rho_{AB}$$

$$\rho_{B|A} = \rho_{A}^{-1/2} \rho_{AB} \rho_{A}^{-1/2}$$

$$\rho_{B} = \operatorname{Tr}_{A}(\rho_{B|A} \rho_{A})$$

 $ho_{B|A} \ge 0$

$$P(R,S) \qquad \qquad \begin{array}{c} \varrho_{AB} \\ \hline \mathbf{S} \\ \hline \mathbf{R} \end{array} \qquad P(S|R) = P(R,S)/P(R) \qquad \qquad \begin{array}{c} \mathbf{B} \\ \mathbf{B} \\ \mathbf{P}(S|R) = P(R,S)/P(R) \end{array} \qquad \qquad \begin{array}{c} \varrho_{B|A} = \varrho_A^{-1/2} \varrho_{AB} \rho_A^{-1/2} \\ \rho_B = \operatorname{Tr}_A(\varrho_{B|A} \rho_A) \end{array}$$

 $arrho_{B|A}^{T_A} \geq 0$

Comparing causal and acausal correlations in Quantum Mechanics



	$\widehat{\boldsymbol{Q}}_{A'} \widehat{\boldsymbol{Q}}_{B}$	$\widehat{P}_{A}, \widehat{P}_{B}$
EPR>	С	А
(I⊗U _{inv}) EPR⟩	А	С

	$\widehat{\boldsymbol{Q}}_{A'} \widehat{\boldsymbol{Q}}_{B}$	$\widehat{P}_{A}, \widehat{P}_{B}$
id	С	С
inv	А	А

Comparing causal and acausal correlations in Epistemically Restricted Liouville mechanics

$$P_{id}(q_B, p_B | q_A, p_A)$$

 $\propto \delta(q_A - q_B)\delta(p_A - p_B)$



 $P_{\text{EPR}}(q_A, p_A, q_B, p_B) \propto \delta(q_A - q_B)\delta(p_A + p_B)$

$$Q_B - Q_A = 0$$
$$P_B + P_A = 0$$

	Q_A, Q_B	P _A ,PB
P _{EPR}	C	А
P _{EPR-inv}	A	С

	Q_A, QB	P _A ,PB
P _{id}	С	С
P _{inv}	А	А

Comparing causal and acausal correlations in Epistemically Restricted Liouville mechanics

 $P_{\text{EPR}}(q_B, p_B | q_A, p_A)$ $\propto \delta(q_A - q_B)\delta(p_A + p_B)$



 $\begin{array}{l} P_{\mathrm{id}}(q_B,p_B|q_A,p_A) \propto \delta(q_A-q_B)\delta(p_A-p_B) \\ Q_B-Q_A=0 \\ P_B-P_A=0 \end{array} \text{ Not allowed!} \end{array}$

	Q_A, Q_B	P_A, PB
P _{EPR}	С	А
P _{EPR-inv}	А	С

	Q_A, QB	P _A ,PB
P _{id}	С	С
P _{inv}	А	А

 $P(S|R) \in \Lambda_R(L_{\text{restricted}}) \qquad \qquad \varrho_{B|A}^{T_A} \ge 0$

 $P(S|R) \in \Lambda_R(L_{\text{restricted}})$

 $\tau_{B|A} \in T_A(O_{\text{restricted}})$

Quantum steering and quantum collapse in the conditional states formalism

Quantum steering

Measure: $\rho_{X|A} \leftrightarrow \{E_x^A\} \qquad \text{Infer:} \quad \rho_B \rightarrow \rho_{B|X=x}$ Learn: X = x $A \rightarrow --- B$ $\rho_{B|X} = \rho_X^{-1/2} \rho_B^{1/2} \rho_{X|B} \rho_B^{1/2} \rho_X^{-1/2}$ where $\rho_{X|B} = \text{Tr}_A(\rho_{X|A} \rho_{A|B})$ $\rho_{A|B} = \rho_B^{-1/2} \rho_{AB} \rho_B^{-1/2}$

$$\begin{split} \rho_{B|X=x} &= \frac{\rho_B^{1/2} E_x^B \rho_B^{1/2}}{\text{Tr}_B(E_x^B \rho_B)}\\ \text{where} \quad E_x^B &= \mathfrak{E}^{\dagger}_{A \to B}(E_x^A) \end{split}$$

Fuchs, quant-ph/0205039



Collapse rule:

$$\rho_B \to \frac{(E_x^B)^{1/2} \rho_B (E_x^B)^{1/2}}{\text{Tr}_B (E_x^B \rho_B)}$$

Pure Bayesian conditioning:

$$\rho_B \rightarrow \frac{\rho_B^{1/2} E_x^B \rho_B^{1/2}}{\text{Tr}_B(E_x^B \rho_B)}$$



Given: $\varrho_{A'B'|AB} \leftrightarrow U_{AB \rightarrow A'B'}$

A' gets info about BB' gets info about A(i.e. mmt is informative)→(B to B' is not identity channel)

no information gain without disturbance

Measure:

$$\rho_{X|A'} \leftrightarrow \{\Pi_x^{A'}\}$$
 $A' B' \rho_{B'} = \sum_x P(X = x) \frac{(E_x^B)^{1/2} \rho_B(E_x^B)^{1/2}}{\text{Tr}_B(E_x^B \rho_B)}$
 P_B





No local explanation of Bell inequality violations



Further work

- Conditional independence, sufficient statistics
- Retrodiction, Pre and post selection, General inference
- State compatibility and state pooling

References

- Bartlett, Rudolph and Spekkens, Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction, arXiv:1111:5057
- Leifer, *Quantum Dynamics as an anolog of conditional probability,* PRA 74, 042310 (2006)
- Leifer and Spekkens, Formulating Quantum Theory as a Causally Neutral Theory of Bayesian Inference, arXiv:1107.5849
- Wood and Spekkens, *The lesson of causal discovery algorithms for quantum correlations*, arXiv:1208.4119

	Conventional expression	In terms of conditional states
Born's rule	$\forall y : P(Y = y) = \operatorname{Tr}_A(E_y^A \rho_A)$	$\rho_Y = \operatorname{Tr}_A(\rho_{Y A}\rho_A)$
Ensemble averaging	$\rho_A = \sum_x P(X = x) \rho_x^A$	$\rho_A = \operatorname{Tr}_X(\rho_{A X}\rho_X)$
Action of quantu channel	m $ ho_B = \mathcal{E}^{A o B}(ho_A)$	$\rho_B = \operatorname{Tr}_A(\varrho_{B A}\rho_A)$
Composition of channels	$\mathcal{E}^{A \to C} = \mathcal{E}^{B \to C} \circ \mathcal{E}^{A \to B}$	$\varrho_{C A} = \operatorname{Tr}_B(\varrho_{C B}\varrho_{B A})$
State update	$\forall y : P(Y = y)\rho_y^B = \mathcal{E}_y^{A \to B}(\rho_A)$	$\rho_{YB} = \operatorname{Tr}_A(\varrho_{YB A}\rho_A)$

An idea for achieving realism in quantum theory

Some modification of our classical theory of Bayesian inference

Some modification of our classical theory of reality

Quantum retrodiction

=



Infer: $\rho_B \to \rho_B |_{X=x}$ $\rho_B|_X = \operatorname{Tr}_A(\varrho_B|_A \rho_A|_X)$ $\varrho_B|_A = \varrho_A|_B * (\rho_B \rho_A^{-1})$ $\rho_A|_X = \rho_X|_A * (\rho_A \rho_X^{-1})$ $\rho_B|_{X=x} = \frac{\rho_B^{1/2} \mathcal{E}^{\dagger}{}_{A \to B} (E_x^A) \rho_B^{1/2}}{\operatorname{Tr}_A (E_x^A \mathcal{E}_B \to A(\rho_B))}$

Generalizes Barnett, Pegg & Jeffers, J. Mod. Opt. 47:1779 (2000).

Time symmetry:

Set of possible **predictive** inferences Set of possible **retrodictive** inferences

The Jamiolkowski isomorphism

$$\tau_{B|A} = (\Phi_{A' \to B} \otimes \mathrm{id}_A) (\sum_{j,k} |j\rangle \langle k|_{A'} \otimes |k\rangle \langle j|_A)$$

 $\Phi_{A' \to B}$ is trace-preserving $\leftrightarrow \operatorname{Tr}_B \tau_{B|A} = I_A$

This implies:
$$\Phi_{A
ightarrow B}(
ho_A) = {
m Tr}_A(au_{B|A}
ho_A)$$

Proof:
$$\operatorname{Tr}_{A}(\tau_{B|A}\rho_{A}) = \Phi_{A' \to B}(\sum_{j,k} |j\rangle \langle k|_{A'} \langle j|\rho_{A}|k\rangle)$$

= $\Phi_{A \to B}(\sum_{j,k} |j\rangle \langle j|_{A}\rho_{A}|k\rangle \langle k|_{A})$
= $\Phi_{A \to B}(\rho_{A})$ QED

The Jamiolkowski isomorphism

$$\tau_{B|A} = (\Phi_{A' \to B} \otimes \mathrm{id}_A)(\sum_{j,k} |j\rangle \langle k|_{A'} \otimes |k\rangle \langle j|_A)$$

 $\Phi_{A' \to B}$ is trace-preserving $\leftrightarrow \operatorname{Tr}_B \tau_{B|A} = I_A$

This implies:
$$\Phi_{A
ightarrow B}(
ho_A) = {
m Tr}_A(au_{B|A}
ho_A)$$

$$\begin{aligned} \tau_{B|A} &= (\Phi_{A' \to B} \otimes \mathrm{id}_A) (\sum_{j,k} |j\rangle \langle k|_{A'} \otimes |j\rangle \langle k|_A)^{T_A} \\ &= (\Phi_{A' \to B} \otimes \mathrm{id}_A) (d_A |\Psi^+\rangle_{A'A} \langle \Psi^+|)^{T_A} \\ \Phi_{A \to B} \quad \text{is CP} \quad \leftrightarrow \quad \tau_{B|A}^{T_A} \ge 0 \end{aligned}$$

 $\tau_{B|A} = \left(\left[\Phi_{A' \to B} \circ T_{A'} \right] \otimes \operatorname{id}_A \right) \left(d_A |\Psi^+\rangle_{A'A} \langle \Psi^+ | \right)$ $\Phi_{A \to B} \circ T_A \text{ is CP } \leftrightarrow \tau_{B|A} \ge 0$