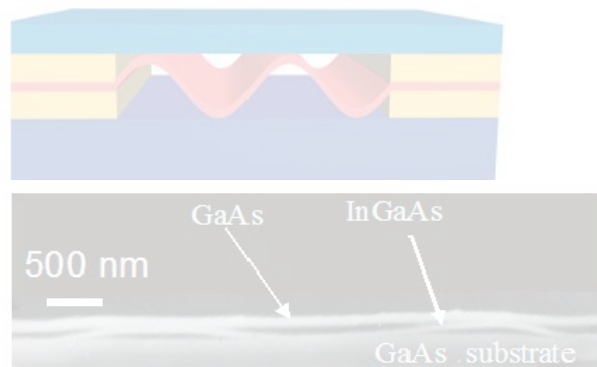


Geometry-Property Relationship in Condensed Matter Physics

Hiroyuki Shima

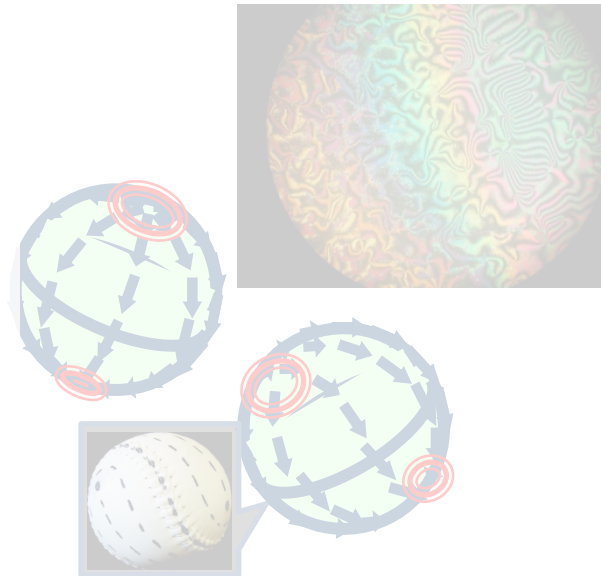
Dept. of Environmental Sciences
University of Yamanashi, Japan

Quantum mechanics on Curved surfaces



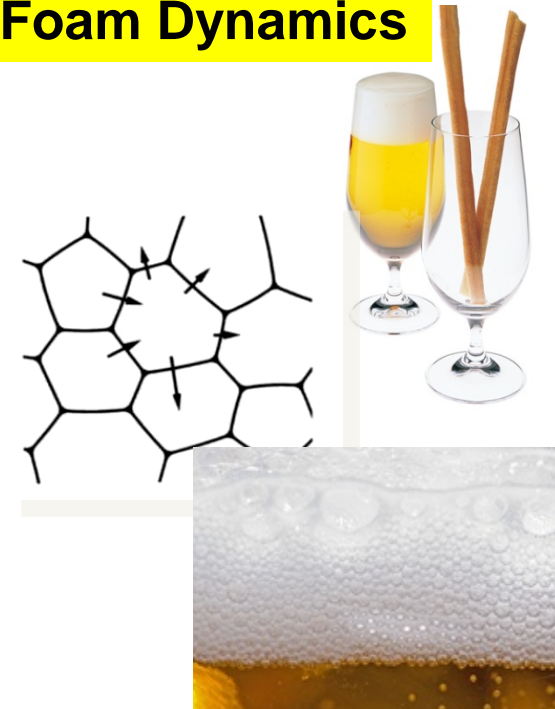
EPL **96** 27011 (2011)
J. Phys. Cond. Mat. **22**, 075301 (2010)
Phys. Rev. B **79**, 201401(R) (2009)
Phys. Rev. B **79**, 235407 (2009)

Liquid crystal on Curved substrates



J. Phys. Soc. Jpn. **79**, 074607 (2010).

Foam Dynamics



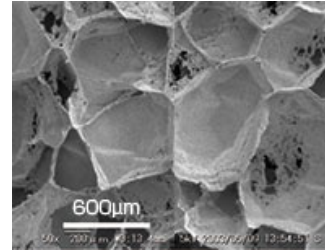
J. Phys. Soc. Jpn. **79**, 074601 (2010).

1. Foam

Shaving cream



Foamed cement



Bread



Bamboo charcoal



Thighbone



Common foam = a dispersion of gas bubble in a liquid

Solid foam = a dispersion of gas bubble in a gel phase (of solid)

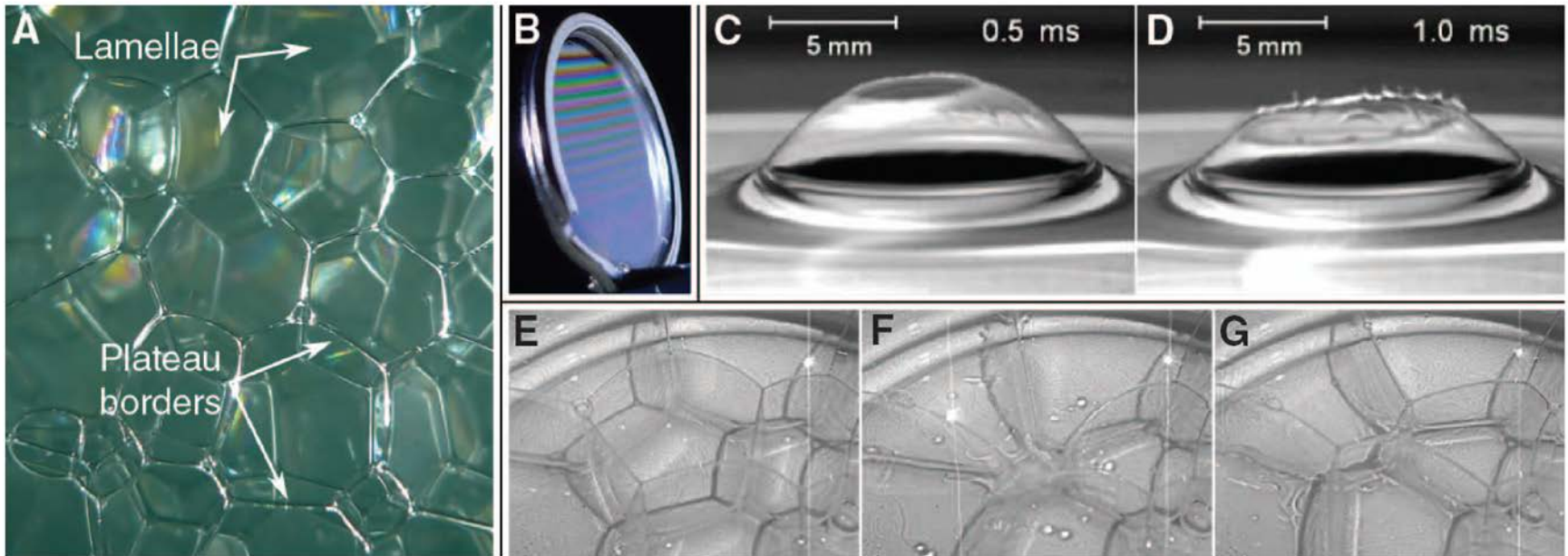
can be made by...

- Gas injection (into soapy water, or through chemical reaction)
- Mechanically beating or stirring

1. Foam

= Treasure box of “physics” and “mathematics”

Saye & Sethian, Science **340** (2013) 720.



- What structure is the **stable** energetically? [surface tension .vs. gas pressure]
- How **drainage** (i.e., downward discharge) of liquid work on the structural change?
- Which cell does **grow**? Which **shrink**?

1. Foam

Multi-scale simulation:

Saye & Sethian, Science 3340 (2013) 720.

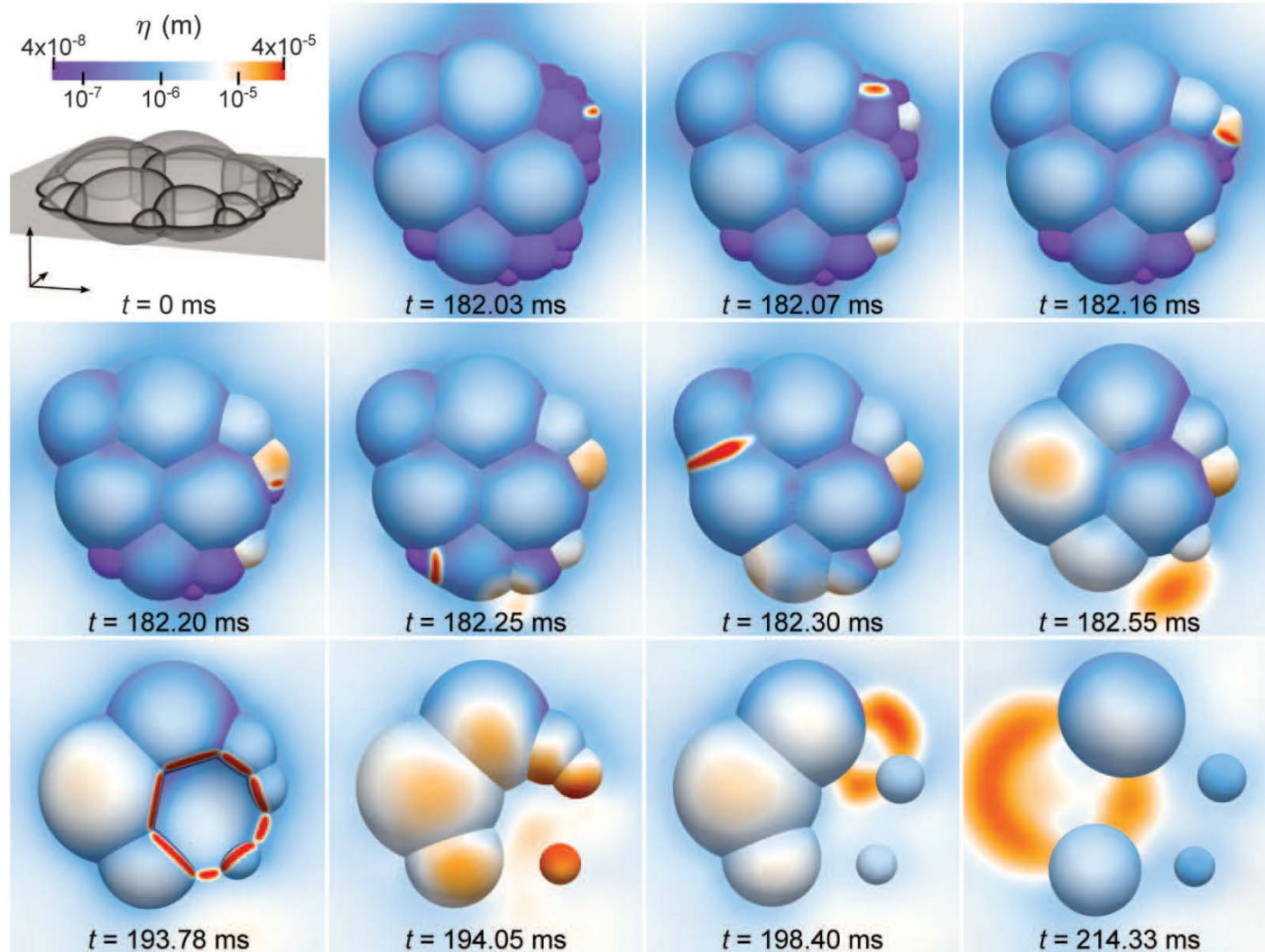
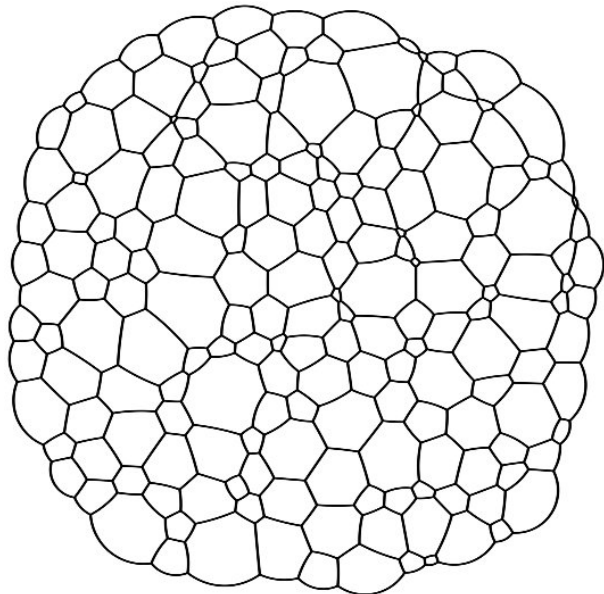
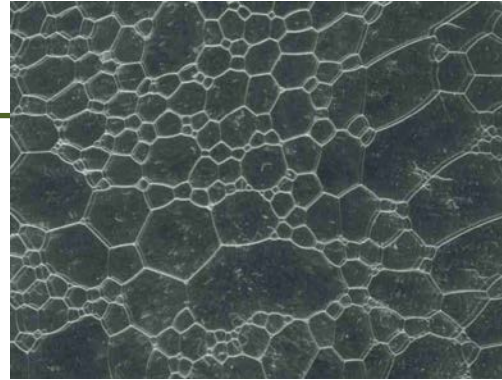


Fig. 3. Results of the coupled multiscale model for a cluster of bubbles attached to a membrane.

2. Foam confined into 2D

Sandwich foam into the gap of two parallel glass plates,
and then look it from above.



We will observe that...

smaller bubbles dissolve/shrink/disappear;

bigger bubbles grow in size;

some liquid interfaces rupture...

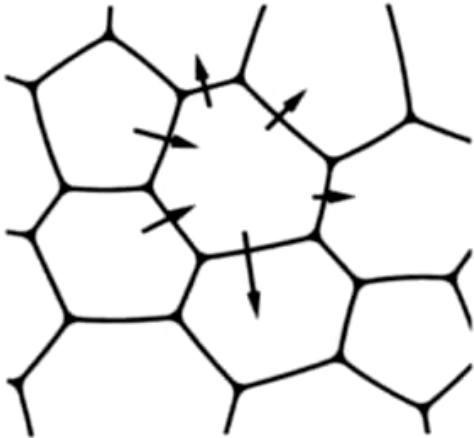
→ **Coarsening of foam**

= Time evolution of a polygons' network

2. Structural instability

(1) Coarsening of air bubbles

... results from internal gas diffusion from a bubble to adjacent another .



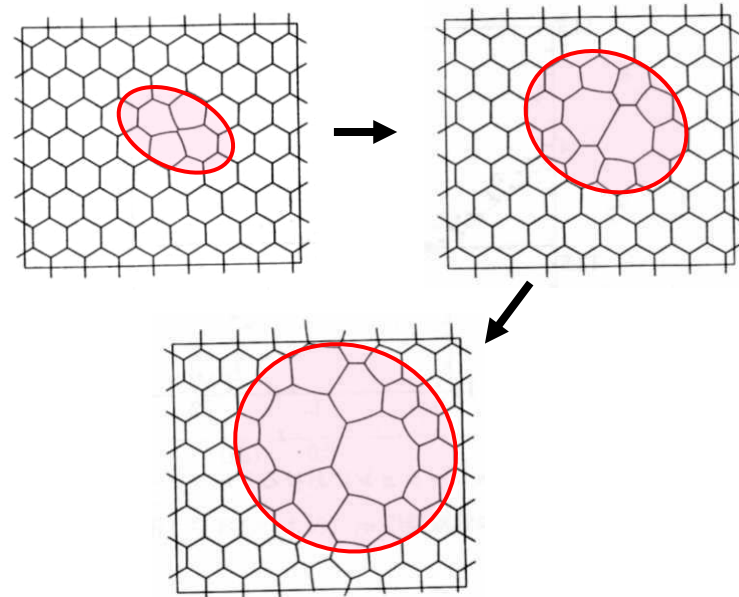
von Neumann's law

$$\text{Flat 2D: } \frac{dS_n}{dt} = \frac{2\pi}{3} \gamma \sigma (n - 6)$$

J. von Neumann, in *Metal Interfaces* (1952).

(2) Rupture of liquid films

... results in a merge of neighboring bubbles and topological changes in the structure



Major causes:

- Drainage & evaporation from liquid films
- Inhomogeneity in interfacial activity
- Contamination

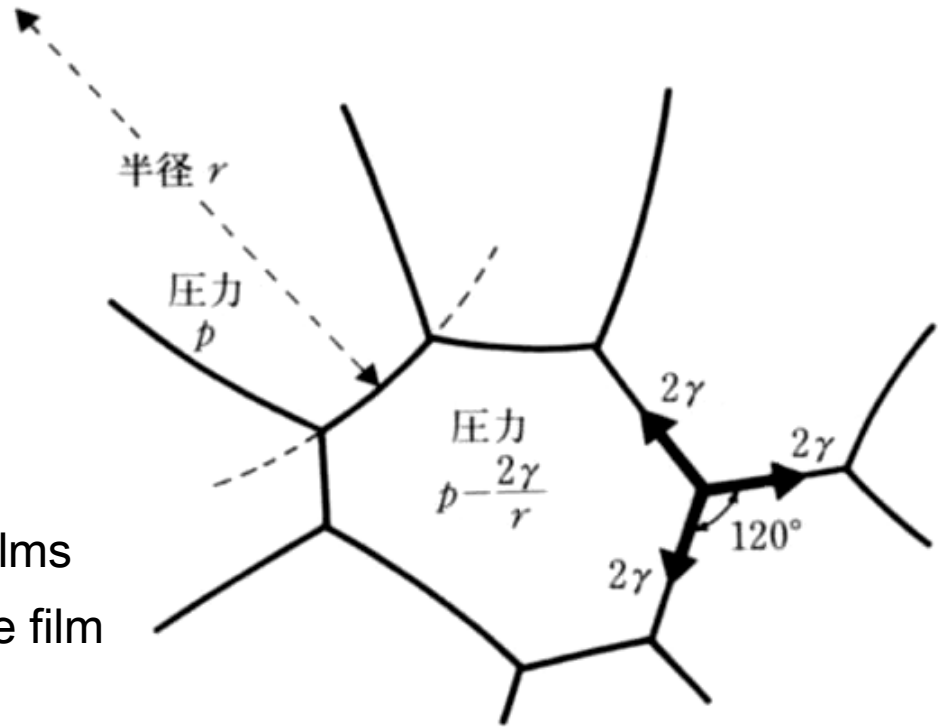
3. Growth of foam

von Neumann relation

$$2D: \frac{dS_n}{dt} = \frac{2\pi}{3} \gamma \sigma (n - 6)$$

σ : surface tension of the liquid films

γ : permeability of gas across the film



(1) Plateau's law : Three films meet along an edge at dihedral angle of 120°

(2) Fick's law $\frac{dS_n}{dt} = -\gamma \sum_j (p_n - p_j) \ell_j$: Pressure-driven gas diffusion across thin liquid films

(3) Laplace's law $\Delta p_i = \frac{2\sigma}{R_i}$: Relation between the pressure difference and the curvature radius of the film

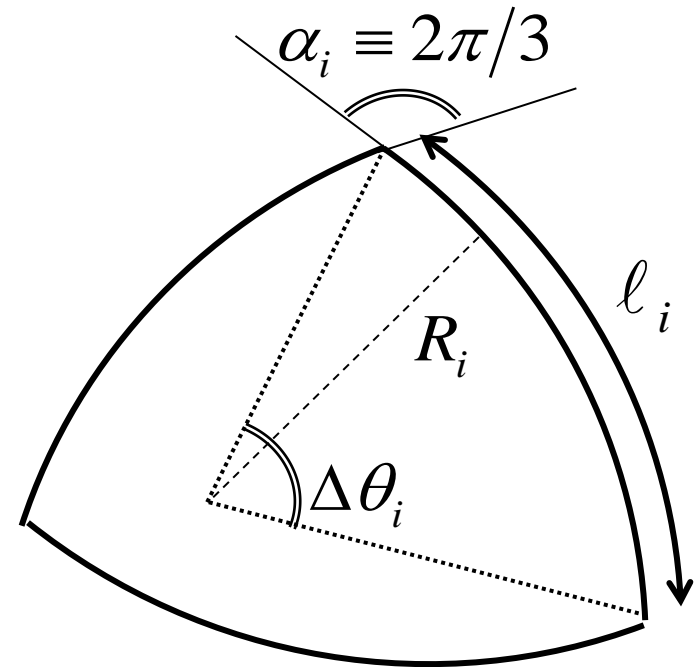
(4) Sum rule of curvature $\sum_{i=1}^n (\pi - \alpha_i) + \sum_{j=1}^n \frac{\ell_j}{R_j} = 2\pi$: Mathematical theorem

4. Sum rule of curvature radius

(1) On a flat plane:

$$\sum_{i=1}^n (\pi - \alpha_i) + \sum_{j=1}^n \frac{\ell_j}{R_j} = 2\pi$$

$$\therefore \left\{ \begin{array}{l} \frac{\pi}{3} \times n + \sum_i \Delta\theta_i = 2\pi \\ \Delta\theta_i = \frac{\ell_i}{R_i} \end{array} \right.$$



(2) On a curved surface with Gaussian curvature K :

$$\sum_{i=1}^n (\pi - \alpha_i) + \sum_{j=1}^n \frac{\ell_j}{R_j} = 2\left(\pi + \iint K dS\right) \quad \dots \text{requires modification of Neumann's law.}$$

5. Foam growth on curved surfaces

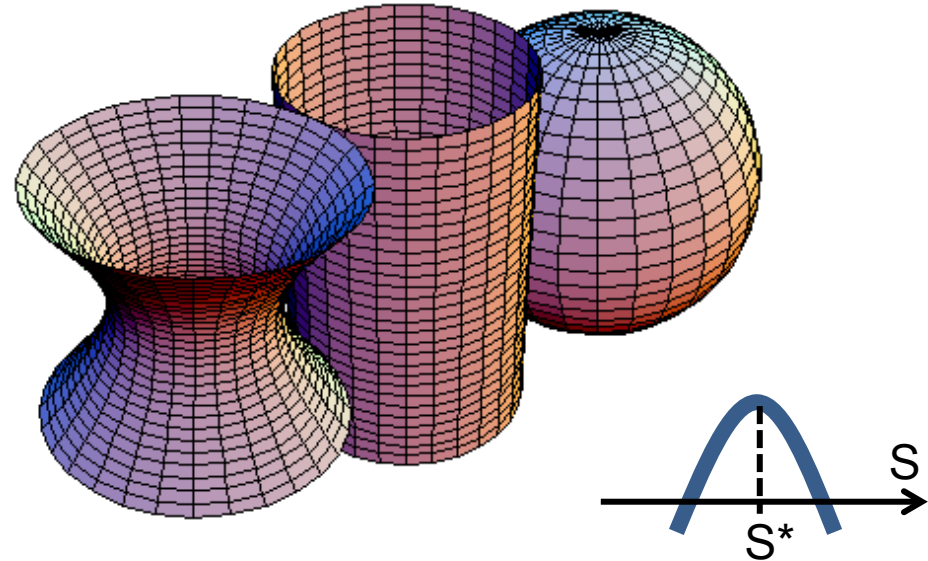
J.E.Avron & D.Levine, Phys.Rev.Lett.**69** (1992) 208.
H.Shima, J.Phys.Soc.Jpn. **79** (2010) 074601.

Neumann's law on curved 2D:

$$\frac{dS}{dt} = \gamma\sigma \left[\frac{\pi}{3} (n - 6) + KS \right]$$

Stationary area of gas babbles:

$$S = S^* \equiv \left| \frac{\pi}{K} \left(2 - \frac{n}{3} \right) \right|$$

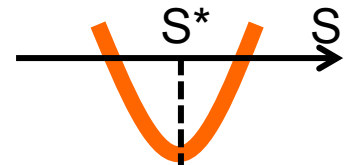


When $K > 0$: S (slightly larger than S^*) makes $dS/dt > 0$

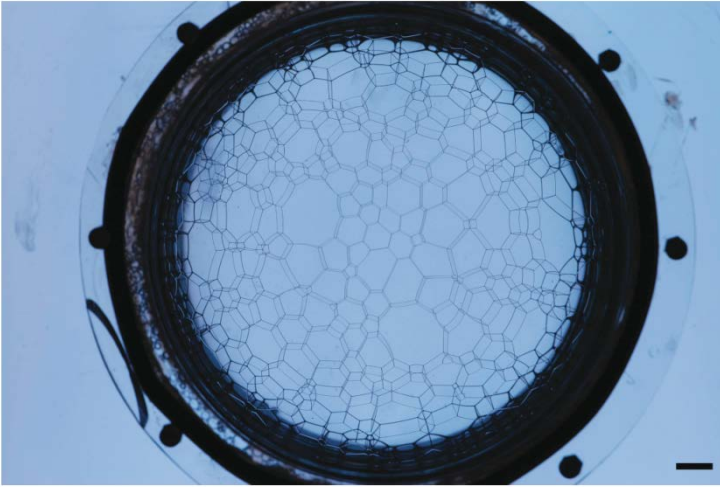
→ All stationary cells are **unstable** (i.e., grow unlimitedly)

When $K < 0$: S (slightly larger than S^*) makes $dS/dt < 0$

→ All stationary cells are **stable**



5. Foam growth on curved surfaces

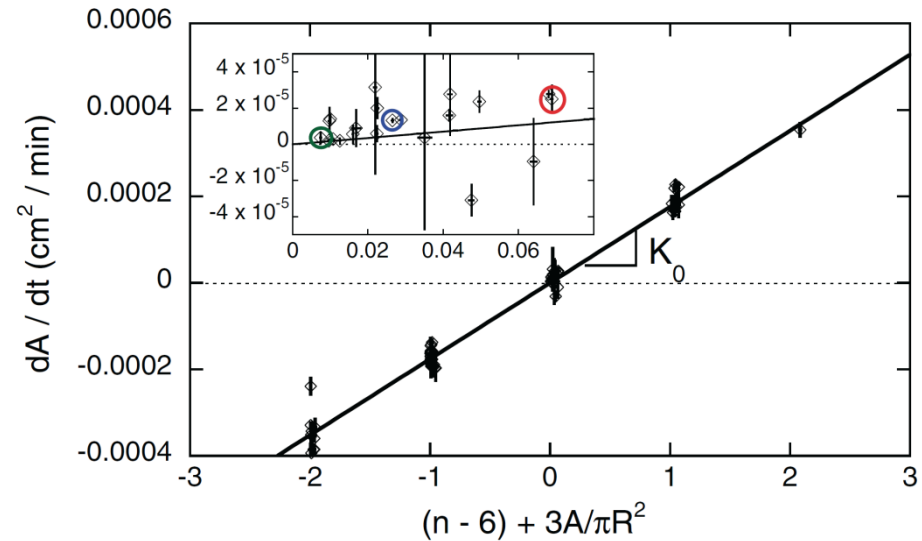
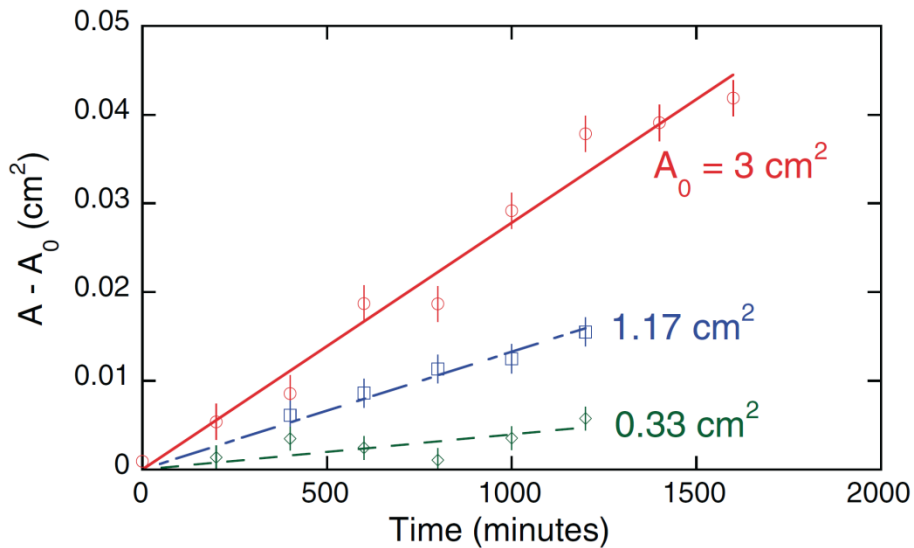


Top view of 2D foam on a hemi-sphere

Roth et al., Phys. Rev. E (2012)

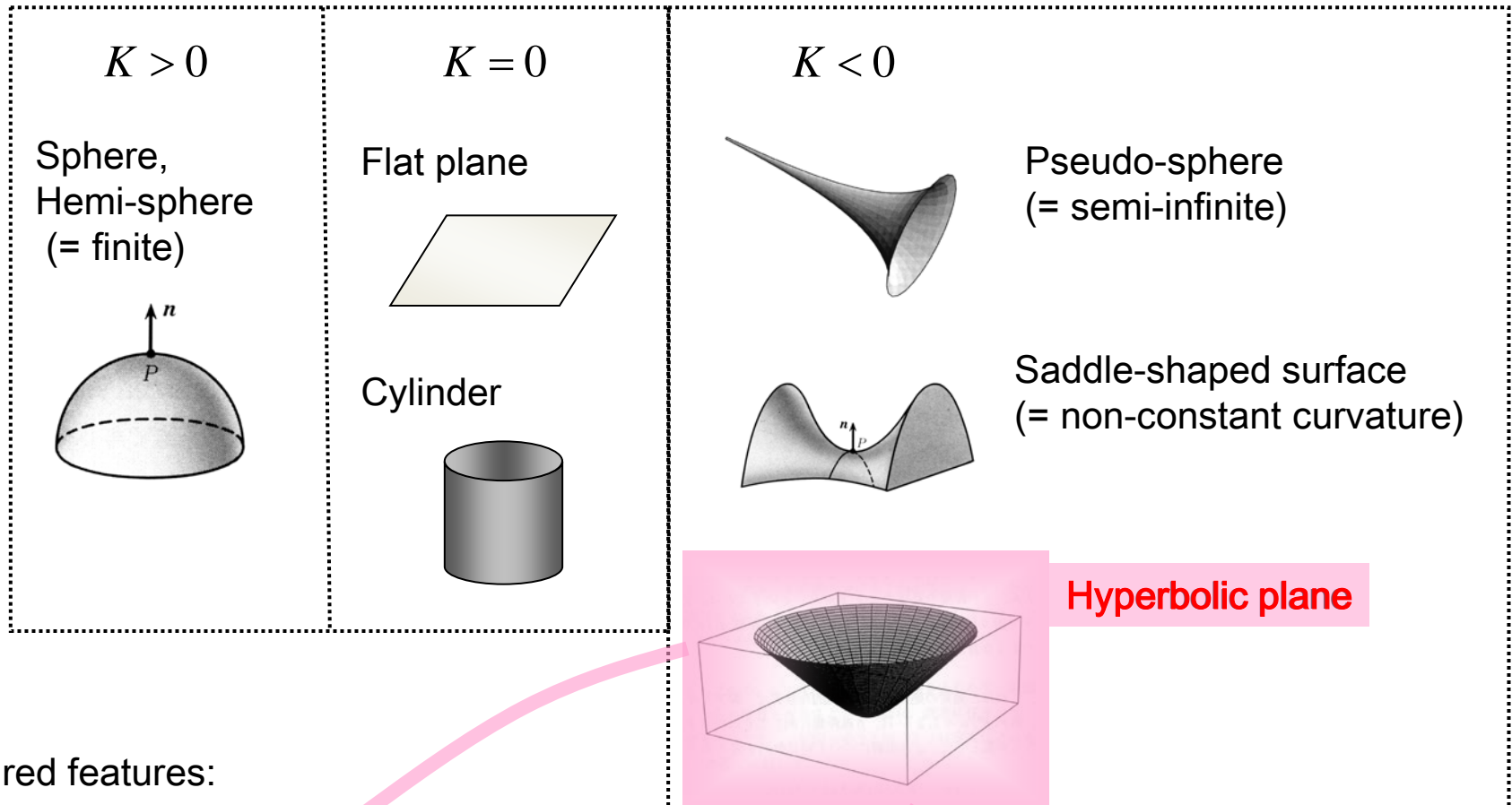
= First experimental test of the curvature effect

$$\frac{dA}{dt} = K_o \left[(n - 6) + \frac{3A}{\pi R^2} \right]$$



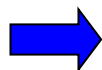
6. On a hyperpolic plane

Classification of curved surfaces:



Desired features:

- (1) an **infinitely-extended** surface;
- (2) possessing a **constant (non-zero) curvature**.

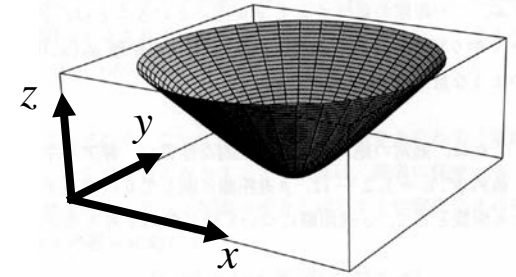


These features enable to extract purely **curvature effects** on foam dynamics.

6. On a hyperpolic plane

Hyperbolic plane

= the locus of points in the **Minkowskian** space whose squared distance from the origin are equal to **-1**



The infinitesimal distance:

$$ds^2 = dx^2 + dy^2 - dz^2$$

The Gaussian curvature:

$$K = -1 \text{ (staying constant)}$$

Spherical surface

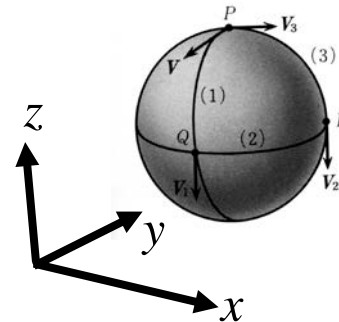
= the locus of points in the **Euclidean** space whose squared distance from the origin are equal to **+1**

The infinitesimal distance:

$$ds^2 = dx^2 + dy^2 + dz^2$$

The Gaussian curvature:

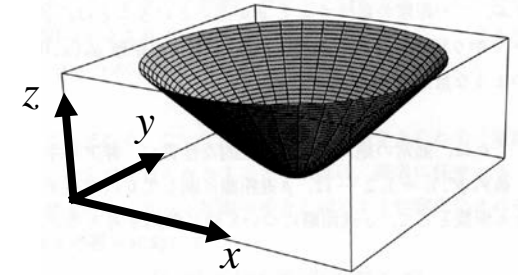
$$K = +1 \text{ (staying constant)}$$



6. On a hyperbolic plane

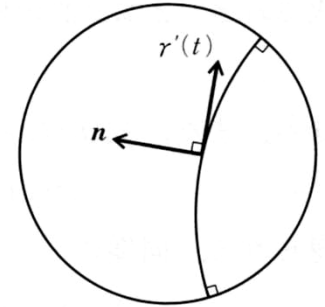
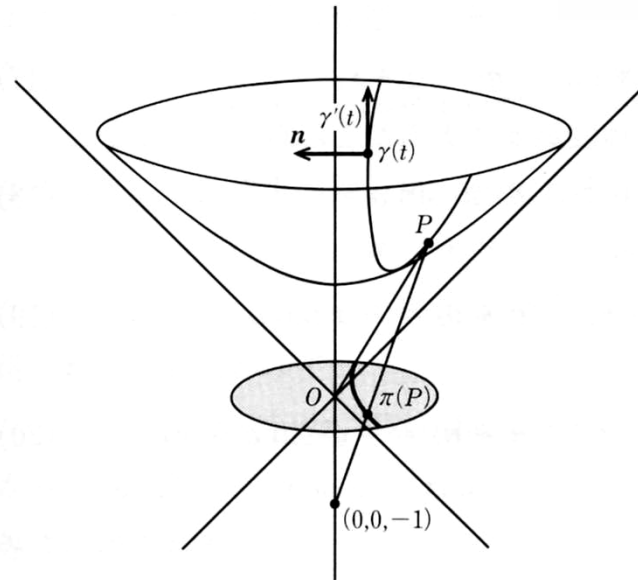
Hyperbolic plane

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The infinitesimal distance:
 $ds^2 = dx^2 + dy^2 - dz^2$

The Gaussian curvature:
 $K = -1$ (staying constant)



Main advantages :

- a simply-connected **infinite** surface
- having a **constant negative** Gaussian curvature
- allowing the establishing of **an infinite variety** of regular lattices
- being expressed in terms of **Poincaré-disk** representation.

6. On a hyperbolic plane

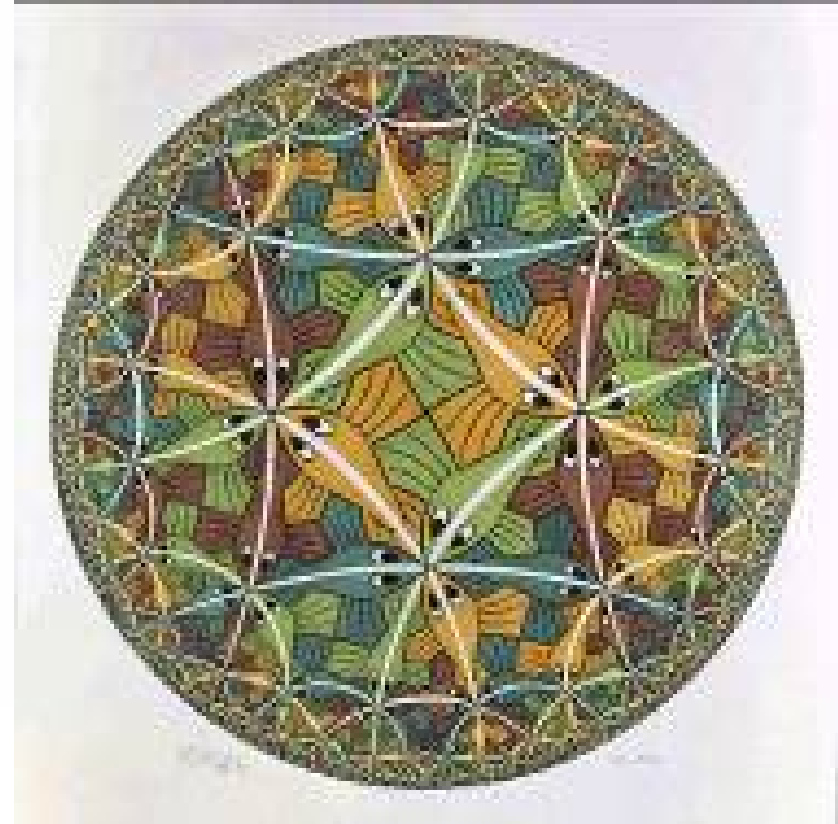
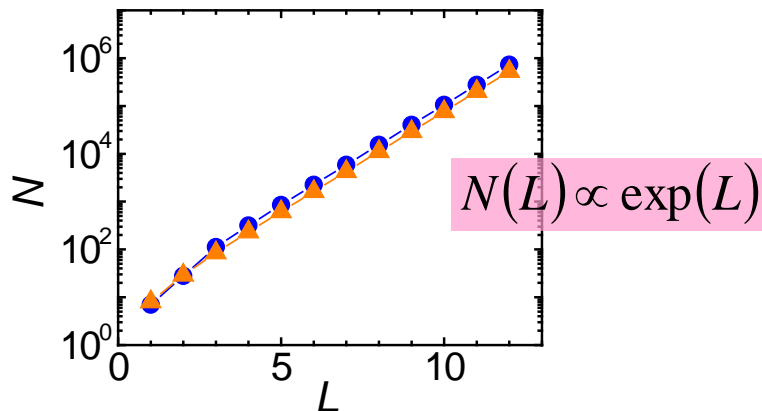
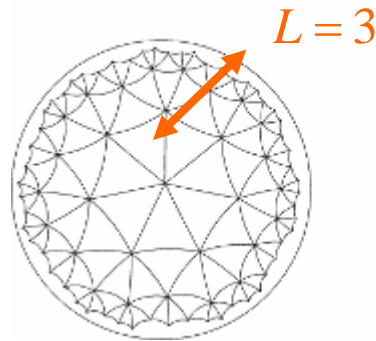
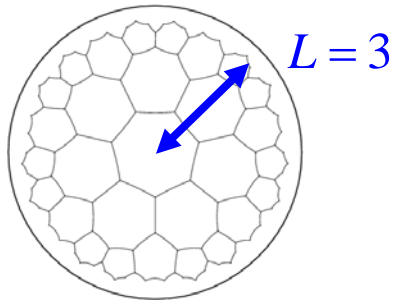
An infinite variety of regular lattices:

Geometric condition: $(p-2)(q-2) > 4$

... At each vertex, q regular p -sided polygons assumes to meet.

Heptagonal
 $[p, q] = [7, 3]$

Triangular
 $[p, q] = [3, 7]$



“Circle Limit III” (M.C.Escher, 1959)
www.mcescher.com/Gallery/gallery.html

6. On a hyperpolic plane

P. Peczak et al., Phys. Rev. E **48** (1993) 4470.

→ von-Neumann law is true for “ $K < 0$ ”

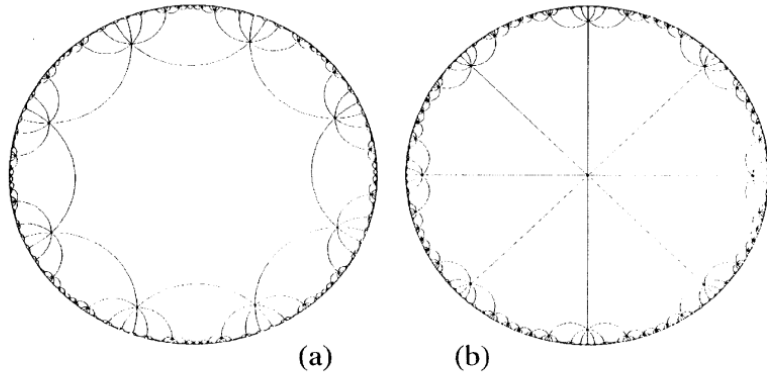
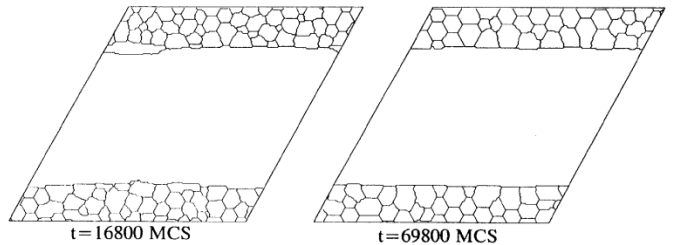
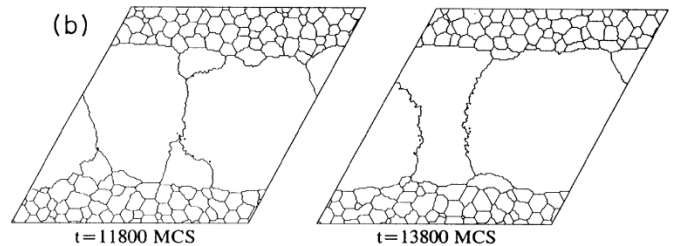
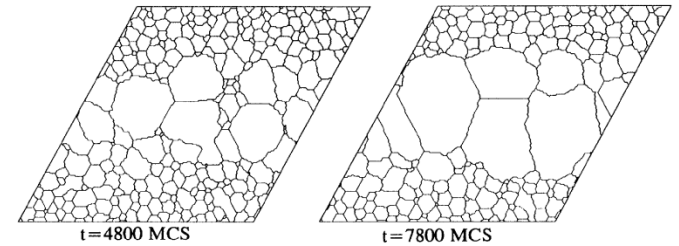
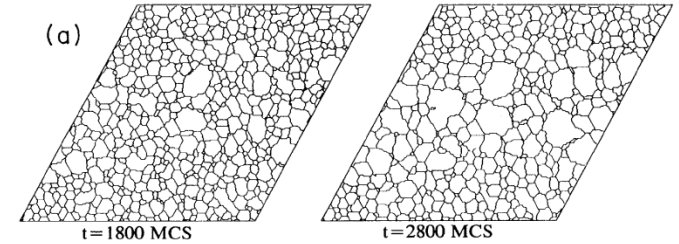
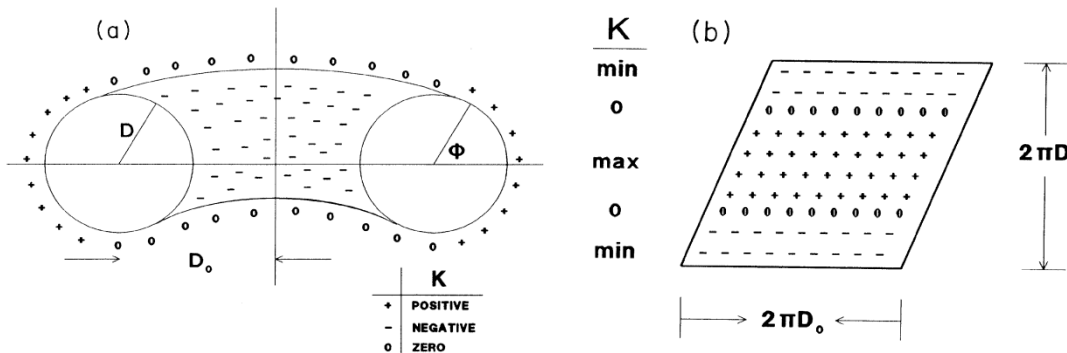


FIG. 12. The octagonal network obtained on the pseudosphere by successive application of fundamental boosts g_0, g_1, g_2, g_3 and their inverses to (a) a fundamental domain and (b) a fundamental vertex (see the text). A tessellation is shown in the Poincaré disc (compare Ref. [52]).



7. Growth in 3D foam

Fick's law

$$\frac{dS}{dt} = \gamma \times \ell_i \times \Delta p_i$$

$$\frac{dV}{dt} = \int_{Face} \gamma_{3D} \Delta p_i dA$$

Laplace's law

$$\Delta p_i = \frac{2\sigma}{R_i}$$

$$\Delta p_i = \sigma \left(\frac{1}{R_{\max}} + \frac{1}{R_{\min}} \right) = 2\sigma H$$

Sum rule of curvature

$$\sum_i \frac{\ell_i}{R_i} = \frac{\pi}{3} (6 - n)$$

None!

... even though, for Gaussian curvature K ,

$$\int_{Face} K dA = 2\pi - n(\pi - \theta_0)$$

$$2D: \frac{dS_n}{dt} = \frac{2\pi}{3} \gamma \sigma (n - 6)$$

$$3D: \frac{dV_n}{dt} = -2\kappa\gamma \int_{Face} H dA = -\frac{2\pi}{3} \kappa\gamma \left[6L(\mathbf{D}) - \sum_{i=1}^n e_i(\mathbf{D}) \right]$$

R.D. MacPherson *et al.*,
 Nature **446** (2007) 1053

e_i : Length of the i -th edge

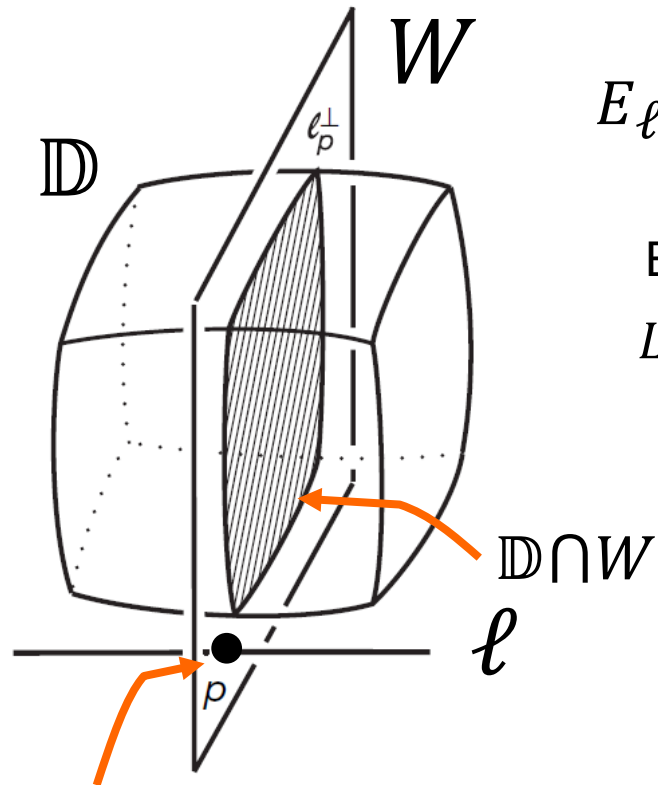
L : Mean width of the 3d domain D

7. Growth in 3D foam

$$3D: \frac{dV_n}{dt} = -2\kappa\gamma \int_{Face} H dA = -\frac{2\pi}{3} \kappa\gamma \left[6L(\mathbb{D}) - \sum_{i=1}^n e_i(\mathbb{D}) \right]$$

R.D. MacPherson *et al.*, Nature **446** (2007) 1053

Cross-section:



p : intersection point

$$E_\ell(\mathbb{D}) = \int_p \chi(\mathbb{D} \cap W) dp$$

Euler characteristics: $\chi=1$ for a no-hole object

$L(\mathbb{D}) = 2E_\ell(\mathbb{D})$ averaged over all possible ℓ

Remarkable fact:

Generalization of Neumann's law
(remaining problem during 50 years)
was resolved by

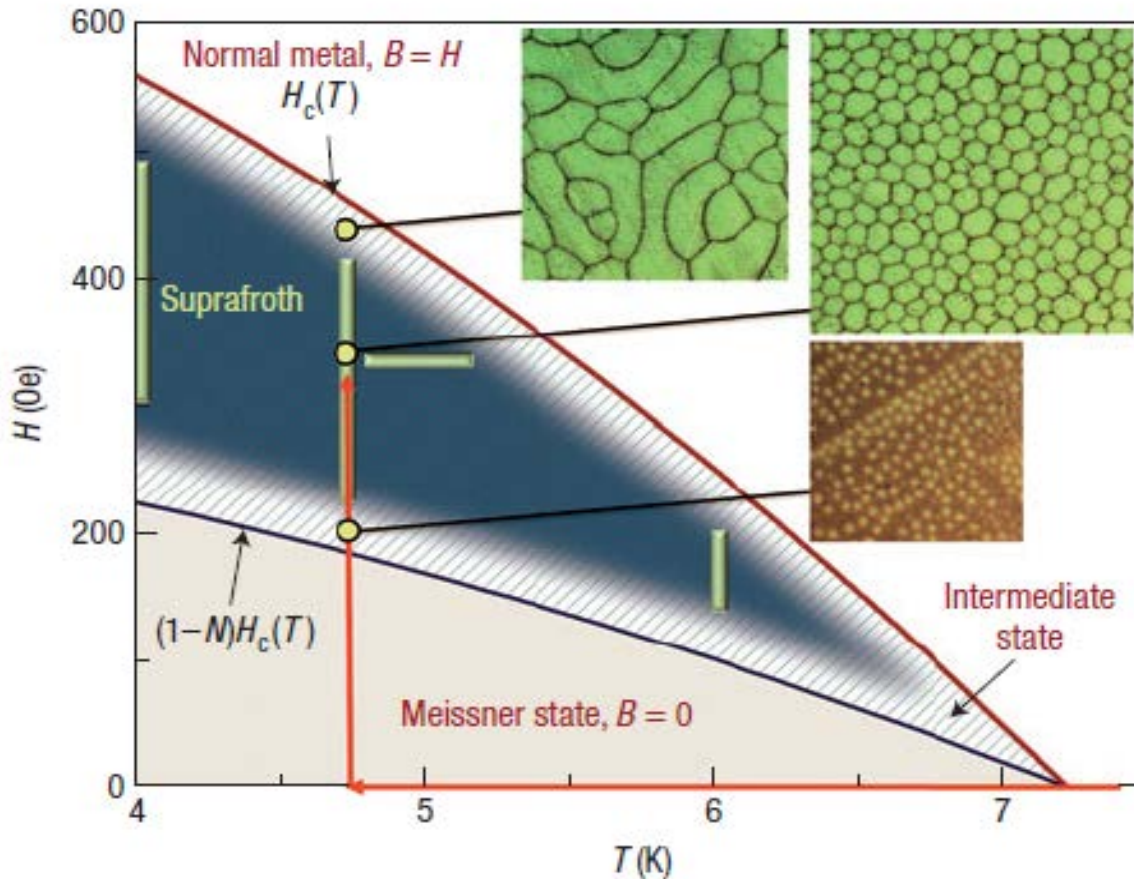
Mathematician: Robert D. MacPherson
and

Physicist: David J. Srolovitz

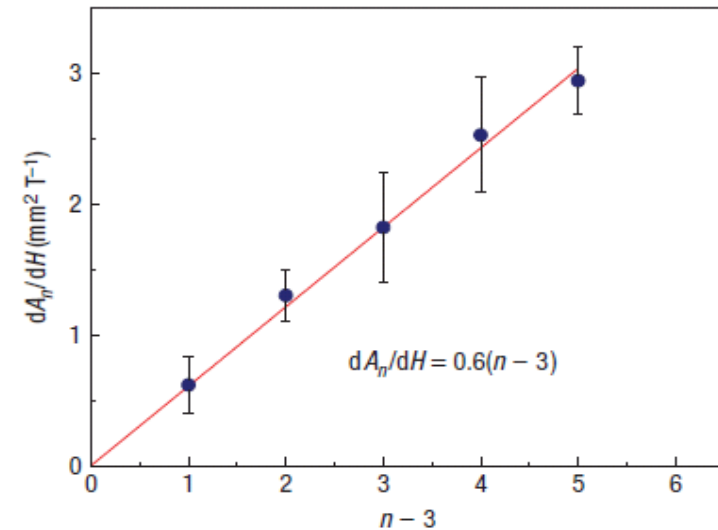
8. Foam in superconducting films

Magnetic field application to a thin Pb disk

➔ Metallic domains (=foam) occur below a critical field strength



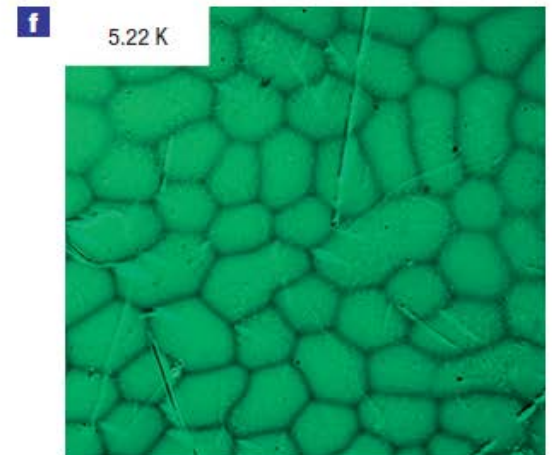
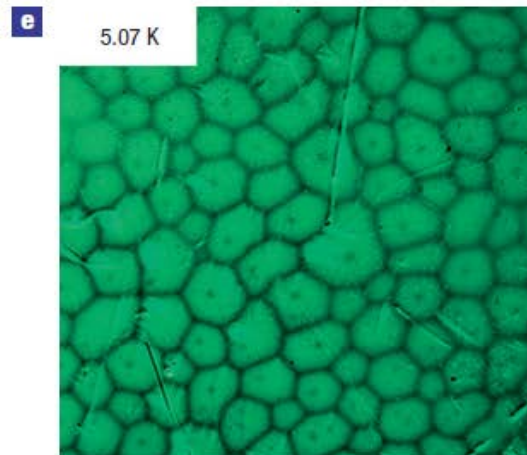
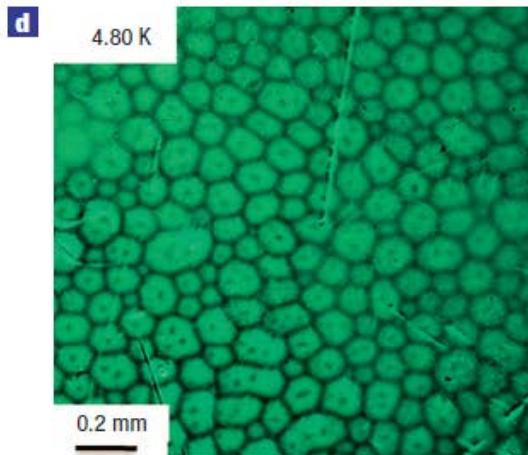
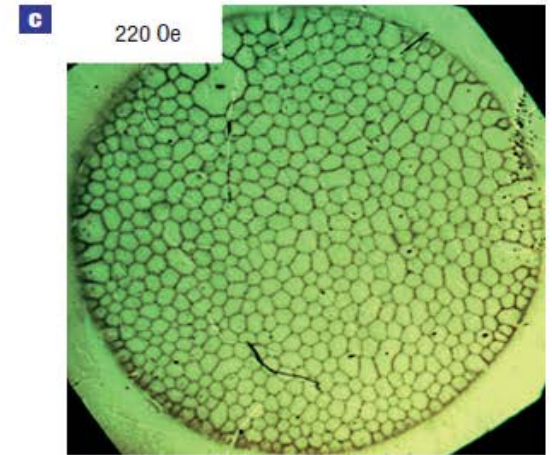
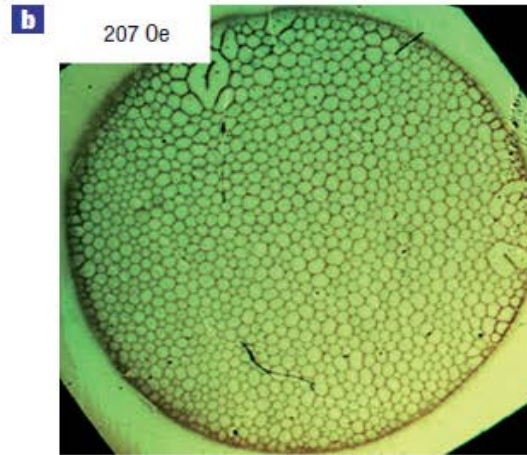
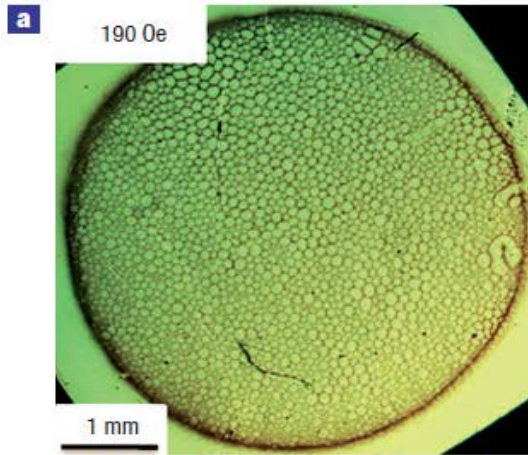
“Suprafoth in type-I superconductors”
R. Prozorov *et al.*,
Nature Physics 4 (2008) 327.



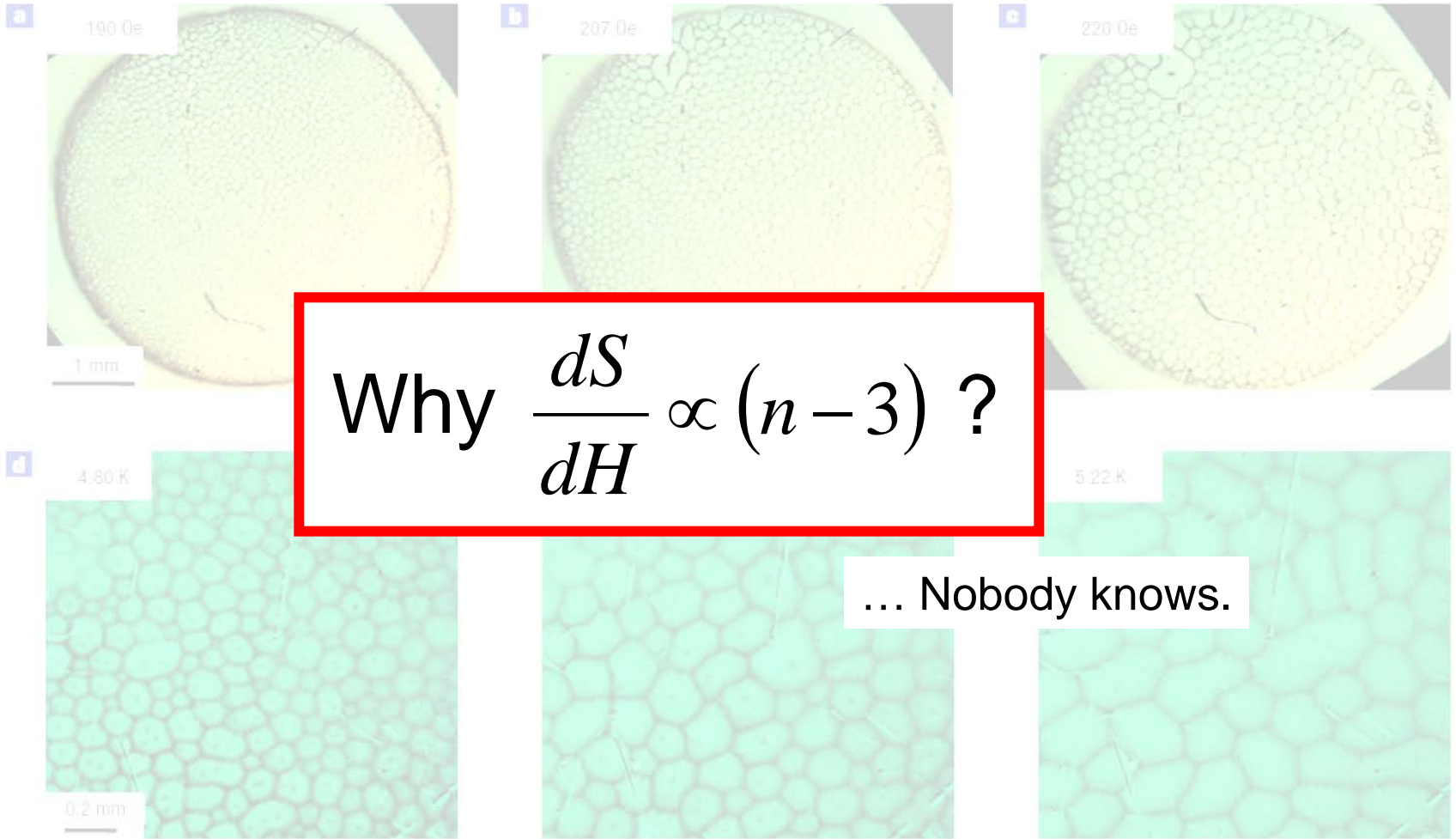
$$\frac{dS}{dH} \propto (n-3)$$

8. Foam in superconducting films

Expository Quantum Lecture Series 2013
(EQualS 2013) at University Putra Malaysia
22-24, November 2013



8. Foam in superconducting films

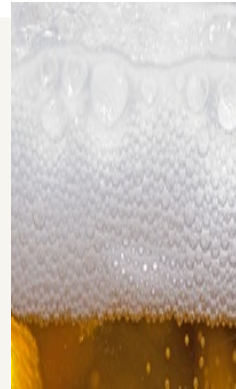
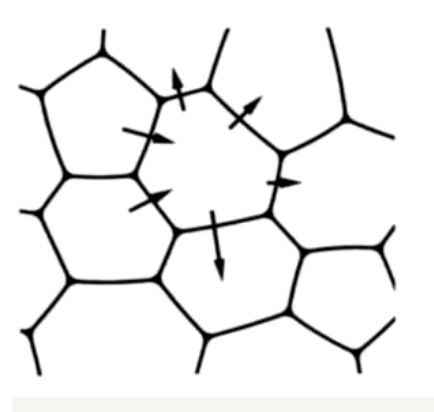


9. Summary

- 2D Foam coarsening

Flat plane:
$$\frac{dS_n}{dt} = \frac{2\pi}{3} \gamma \sigma (n - 6)$$

Curved surface:
$$\frac{dS}{dt} = \gamma \sigma \left[\frac{\pi}{3} (n - 6) + KS \right]$$

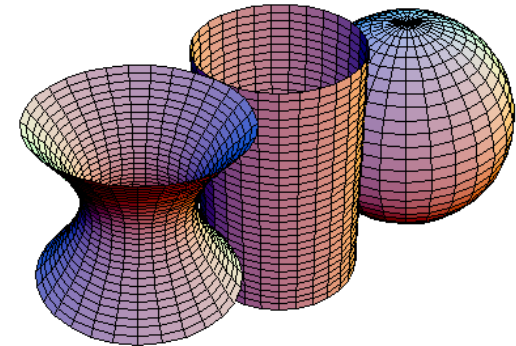


→ When $K > 0$:

All stationary cells are **unstable** (i.e., grow unlimitedly)

→ When $K < 0$:

All stationary cells are **stable**



- 3D Foam coarsening

$$\frac{dV_n}{dt} = -\frac{2\pi}{3} \kappa \gamma \left[6L(\mathbf{D}) - \sum_{i=1}^n e_i(\mathbf{D}) \right]$$

