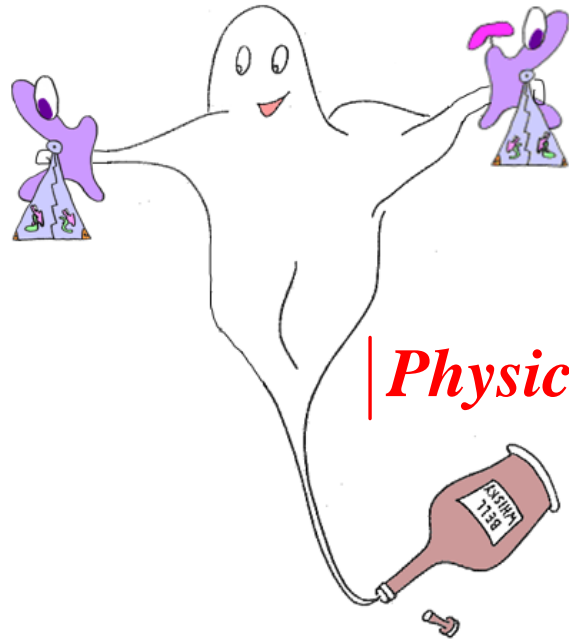


Nonlocality, Entanglement and Decoherence in High Energy Physics

Spooky action at distance
also for neutral kaons?



by
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Institute for Theoretical Physics
University of Vienna
Austria

$$| \textit{Physics} \rangle = \alpha | \textit{Particle Physics} \rangle + \beta | \textit{Quantum Theory} \rangle$$

experimental ↔ phenomenological ↔ conceptual ↔ mathematical aspects



Testing QM in High Energy Physics

- Part I: Bell inequalities 1:

A symmetry violation in particle physics related to nonlocality ?!

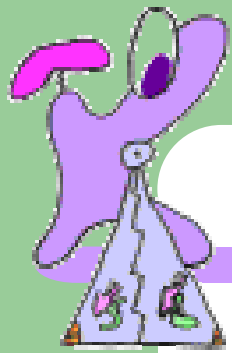
- Part II: The quantum mechanics of neutral kaons (*Exercises*)

- Part III: Bell inequalities 2/ How to describe the decay property?

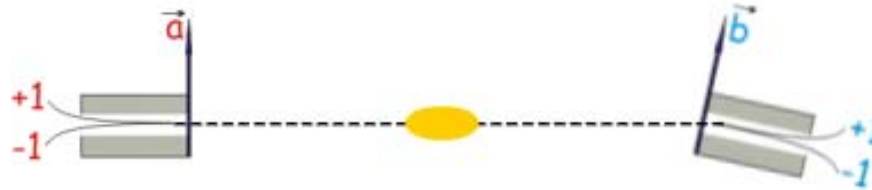
“Erasing the Past and Impacting the Future” by Aharonov & Zubairy



- Part IV: The Kaonic Quantum Eraser/ Decoherence & Measures of Entanglement **or** *Quantum Information* (Sets of entanglement measures of multipartite qudit systems)



The EPR scenario



Antisymmetric Bell state:

$$|\psi^-\rangle = \frac{1}{\sqrt{2}} \left\{ |\uparrow\rangle_l \otimes |\downarrow\rangle_r - |\downarrow\rangle_l \otimes |\uparrow\rangle_r \right\} \quad \dots \text{spin } 1/2$$

$$= \frac{1}{\sqrt{2}} \left\{ |0\rangle_l \otimes |1\rangle_r - |1\rangle_l \otimes |0\rangle_r \right\} \quad \dots \text{qubit}$$

$$= \frac{1}{\sqrt{2}} \left\{ |H\rangle_l \otimes |V\rangle_r - |V\rangle_l \otimes |H\rangle_r \right\} \quad \dots \text{photon}$$

$$= \frac{1}{\sqrt{2}} \left\{ |K^0\rangle_l \otimes |\bar{K}^0\rangle_r - |\bar{K}^0\rangle_l \otimes |K^0\rangle_r \right\} \quad \dots \text{kaon}$$

$$= \frac{1}{\sqrt{2}} \left\{ |B^0\rangle_l \otimes |\bar{B}^0\rangle_r - |\bar{B}^0\rangle_l \otimes |B^0\rangle_r \right\} \quad \dots \text{B-meson}$$

$$= \frac{1}{\sqrt{2}} \left\{ |I\rangle_l \otimes |\uparrow\rangle_r - |II\rangle_l \otimes |\downarrow\rangle_r \right\} \quad \dots \text{single neutron in interferometer}$$

$$= \dots$$



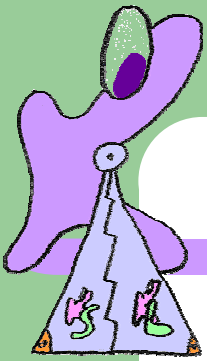
Requirements for tests LRT versus QM

Requirements for a *conclusive* proof of the existence of correlation stronger than those explainable by *locality* and *realism*:

- (1) “*Active*” measurements (opening the possibility for Alice and Bob to choose among alternative setups → free choice)
- (2) “*Use all information*” (test the *whole* ensemble; decay product states are included → this “additional” information should not be ignored)



are not “*only*” loopholes!



Generalized Bell inequality for kaons

$$S_{CHSH}(k_n, k_m, k_{n'}, k_{m'}; t_a, t_b, t_c, t_d)$$

$$= |E(k_n, t_a; k_m, t_b) - E(k_n, t_a; k_{m'}, t_c)| + |E(k_{n'}, t_d; k_m, t_b) + E(k_{n'}, t_d; k_{m'}, t_c)| \leq 2$$

local realistic theories



$$|k_n\rangle = \alpha |K^0\rangle + \beta |\bar{K}^0\rangle \dots \dots \text{quasi-spin}$$

- vary in times
- vary in quasi-spin

QM:

$$S^{\text{Photon}} = 2\sqrt{2} = 2.8$$

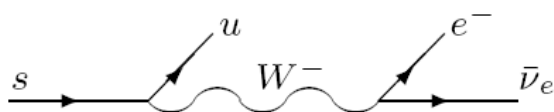
I. Vary in quasi-spin:

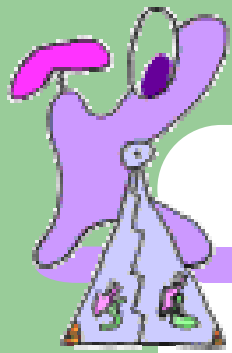
$$\delta = 0$$

CP violation related to
nonlocality!

Leptonic charge asymmetry:

$$\delta = \frac{\Gamma(K_L \rightarrow \pi^- l^+ \nu_l) - \Gamma(K_L \rightarrow \pi^+ l^- \bar{\nu}_l)}{\Gamma(K_L \rightarrow \pi^- l^+ \nu_l) + \Gamma(K_L \rightarrow \pi^+ l^- \bar{\nu}_l)} = (3.27 \pm 0.12) \cdot 10^{-3}$$





Generalized Bell inequality for kaons

$$S_{CHSH}(\bar{K}^0, \bar{K}^0, \bar{K}^0, \bar{K}^0; t_a, t_b, t_c, t_d) \\ = \left| E(\bar{K}^0, t_a; \bar{K}^0, t_b) - E(\bar{K}^0, t_a; \bar{K}^0, t_c) \right| + \left| E(\bar{K}^0, t_d; \bar{K}^0, t_b) + E(\bar{K}^0, t_d; \bar{K}^0, t_c) \right| \leq 2$$

II: vary in times

$$|k_n\rangle = \alpha |K^0\rangle + \beta |\bar{K}^0\rangle \dots \dots \text{quasi-spin}$$

$$|\psi^-\rangle = \frac{1}{\sqrt{2}} \left\{ |K^0\rangle_l \otimes |\bar{K}^0\rangle_r - |\bar{K}^0\rangle_l \otimes |K^0\rangle_r \right\}$$

$$S_{\text{Photon}} = 2\sqrt{2}$$

Violation!

$$S_{\text{Kaon}}\left(\frac{3\pi}{4}, \frac{3\pi}{4}, \frac{\pi}{2}, \mathbf{0}\right) = 1.36$$

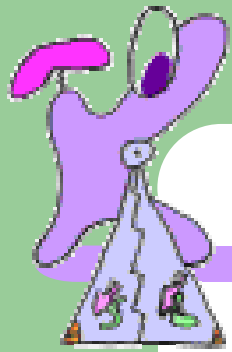
NO violation!

Strangeness oscillation/decay:

$$x = \frac{\Delta m}{\Gamma} \approx \frac{2 \Delta m}{\Gamma_s} \approx 1$$

PROPOSITION:

The CHSH-inequality is violated iff $x > 2$ for kaons.



The dichotomic questions after evolving to $t > 0$?

Photons (or spin $1/2$):

- (1) Are you horizontal or vertical polarized?
- (2) Are you horizontal polarized or not?
- } same!!

(1) Are you a meson K^0 or an antimeson \bar{K}^0 at time t ?

additional information from the decay products
ignored (only dichotomic if conditioned to the
survival of both mesons)

(2) Are you a meson \bar{K}^0 or *not* at time t ?

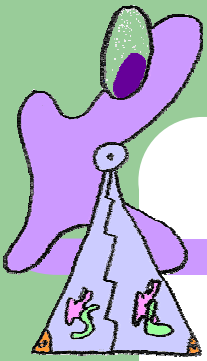
... is dichotomic

Unitary time evolution:

$$U(t, 0) |K_{S/L}\rangle = |K_{S/L}(t)\rangle = e^{-im_{S/L}t - \frac{\Gamma_{S/L}}{2}t} \left(|K_{S/L}\rangle + |\Omega_{S/L}(t)\rangle \right)$$

surviving component
decaying component

$H_{survive} \oplus H_{decay}$



The BI sensitive to strangeness

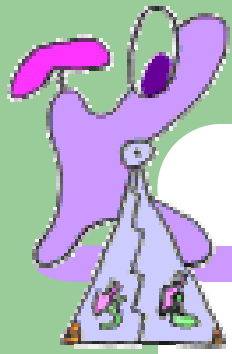
(1) Are you a meson K^0 or an antimeson \bar{K}^0 at time t ?

additional information from the decay products ignored
(only dichotomic if conditioned to the survival of both
mesons)

$$E^{\text{non-unitary}}(t_l; t_r) = P(K^0 t_l; K^0 t_r) + P(\bar{K}^0 t_l; \bar{K}^0 t_r) - P(K^0 t_l; \bar{K}^0 t_r) - P(\bar{K}^0 t_l; K^0 t_r) \cong \\ -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S + \Gamma_L}{2}(t_l + t_r)} \stackrel{\Gamma_L \ll \Gamma_S}{\cong} -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S}{2}(t_l + t_r)}$$

(2) Are you a meson \bar{K}^0 or *not* at time t ?

$$E^{\text{unitary}}(\bar{K}^0, t_l; \bar{K}^0, t_r) = P(Y, t_l; Y, t_r) + P(N, t_l; N, t_r) - P(Y, t_l; N, t_r) - P(N, t_l; Y, t_r) \\ = -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S + \Gamma_L}{2}(t_l + t_r)} \\ + \frac{1}{2}(1 - e^{-\Gamma_S t_l})(1 - e^{-\Gamma_L t_r}) + \frac{1}{2}(1 - e^{-\Gamma_L t_l})(1 - e^{-\Gamma_S t_r}) \\ \stackrel{\Gamma_L \ll \Gamma_S}{\cong} -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S}{2}(t_l + t_r)}$$



The BI sensitive to strangeness

(1) Are you a meson K^0 or an antimeson \bar{K}^0 at time t ?

additional information from the decay products ignored
(only dichotomic if conditioned to the survival of both mesons)

$$E^{non-unitary}(t_l; t_r) = -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S + \Gamma_L}{2}(t_l + t_r)} \stackrel{\Gamma_L \ll \Gamma_S}{\cong} -\cos(\Delta m(t_l - t_r)) \cdot e^{-\frac{\Gamma_S}{2}(t_l + t_r)}$$

$$\sum P = \cosh\left(\frac{\Gamma_L - \Gamma_S}{2}(t_l - t_r)\right) \neq 1$$

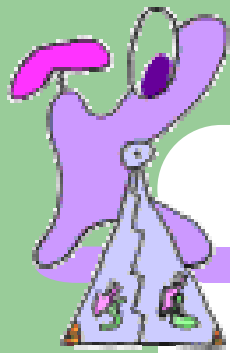
normalize

$$E_{normalized}^{non-unitary}(t_l; t_r) = \frac{-\cos(\Delta m(t_l - t_r))}{\cosh\left(\frac{\Gamma_L - \Gamma_S}{2}(t_l - t_r)\right)}$$

violates *formally*
the BI!

- *strangeness oscillation* = birefringence of photons in asymmetric optical fibers
- *decay* = polarization dependent loss for photons

Gisin, Go
Am.J.Phys. (2001)



The

(1) Are you

E^{nc}

$$\sum P = \cosh\left(\frac{\Gamma_L - \Gamma_S}{2}(t_l - t_r)\right)$$

PhysicsWeb - Mesons violate Bell's inequality - Microsoft Internet Explorer von Lycos Bertelsmann

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Mesons violate Bell's inequality

6 November 2003

The famous Bell's inequality of quantum mechanics has been tested in a high-energy particle physics experiment for the first time. The inequality was violated by three standard deviations in experiments with B mesons at the KEK laboratory in Japan - yet again confirming the predictions of quantum theory (arxiv.org/abs/quant-ph/0310192; J. Mod. Optics to be published). Previously most Bell's inequality experiments have been performed with photons or ions.

Experiments to test Bell's inequality involve measuring the properties of pairs of particles that are space-like separated in the sense of special relativity: in other words, there is no time for a light signal to travel between them within the duration of the experiment. In a typical Bell's inequality experiment the polarizations of a pair of photons are measured as the relative angle between the axes of polarizers making the measurements is varied.

Quantum mechanics predicts that "non-local" correlations can exist between the particles. This means that if one photon is polarized in, say, the vertical direction, the other will always be polarized in the horizontal direction, no matter how far away it is. However, some physicists argue that this cannot be true and that quantum particles must have local values - known as "hidden variables" - that we cannot measure.

Belle experiment

Belle and others showed that it was possible to distinguish between quantum mechanics and these hidden-variable theories in a certain type of experiment that measure a parameter known as S. Put simply, the local theories predict that S will always be less than two, whereas the quantum prediction is $S = 2\sqrt{2}$. When S is greater than two, Bell's inequality is said to be violated.

Apollo Go of the National Central University in Taiwan and co-workers in the Belle collaboration performed the experiment at the KEK B-factory. At this accelerator beams of electrons and positrons are collided to produce pairs of B mesons and their antiparticles, which then decay into lighter particles. The meson pairs behave like photon pairs, but instead of analyzing correlations between directions of polarization, the Belle team study particle-antiparticle correlations using a technique known as "flavour tagging". Go and colleagues calculated that $S = 2.725$, with error bars that mean that the inequality is violated by three standard deviations.

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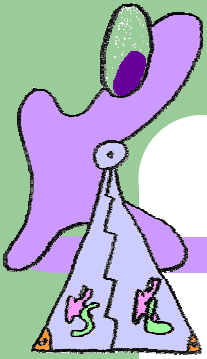
Author

[Peter Rodgers](#)

mayr

$m(t_l - t_r)$

maly



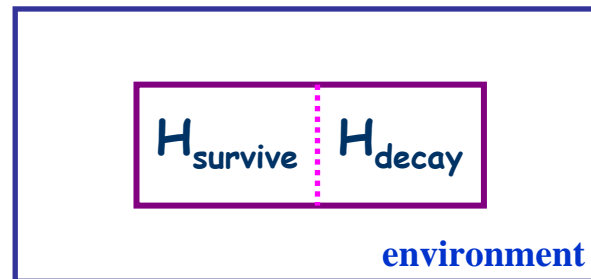
Decay is “kind of decoherence” ?

Kaon in time:

$$|K^0(t)\rangle \cong \frac{1}{2} \left\{ e^{-\frac{\Gamma_S}{2}t - im_S t} |K_S\rangle + e^{-\frac{\Gamma_L}{2}t - im_L t} |K_L\rangle \right\}$$

short-lived state
long-lived state

→ state **not** normalized !!



An open-quantum-system formulation of particle decay

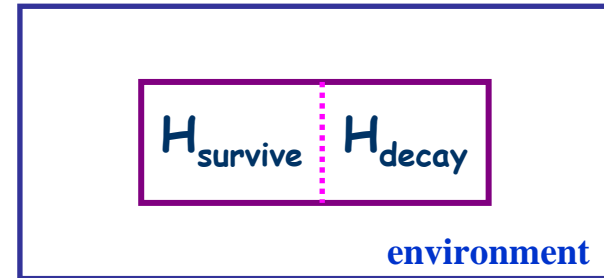
(Bertlmann, Grimus, Hiesmayr; Phys. Rev. A 73, 054101 (2006))

Wigner-Weisskopf approximation is completely positive!

!! particle decay is “kind of decoherence” !!



Decay is “kind of decoherence” ?



Liouville-von Neumann Eq.:

$$\frac{\partial \rho}{\partial t} = -i[H, \rho] - D[\rho]$$

Dissipation/
Decoherence

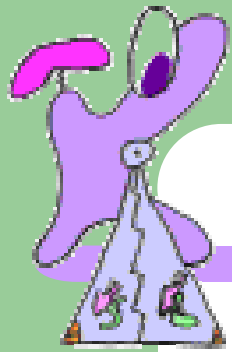
$$D[\rho] = \frac{1}{2} \sum_i \lambda_i (A_i^\dagger A_i \rho + \rho A_i^\dagger A_i - 2A_i \rho A_i^\dagger)$$

1976: Lindblad;
Gorini-Kossakowski-Sudarshan

single kaon described by 4x4 matrix:

$$\rho = \begin{pmatrix} \rho_{ss} & \rho_{sf} \\ \rho_{fs} & \rho_{ff} \end{pmatrix}$$

$s \dots survive; f \dots final / decay$



Open quantum formalism for single kaons

single kaon described by 4x4 matrix:

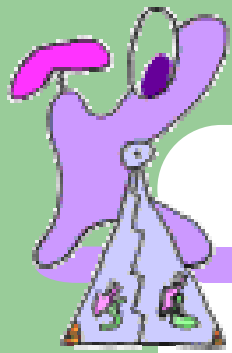
$$\rho = \begin{pmatrix} \rho_{ss} & \rho_{sf} \\ \rho_{fs} & \rho_{ff} \end{pmatrix} \rightarrow \mathbf{0}$$

s ... survive; f ... final / decay



$$\rho(t) = \begin{pmatrix} e^{-\Gamma_s t} \rho_{SS} & e^{-i\Delta m t - \Gamma t} \rho_{SL} & 0 & 0 \\ e^{i\Delta m t - \Gamma t} \rho_{LS} & e^{-\Gamma_L t} \rho_{LL} & 0 & 0 \\ 0 & 0 & (1 - e^{-\Gamma_L t}) \rho_{LL} & X \\ 0 & 0 & X^* & (1 - e^{-\Gamma_s t}) \rho_{SS} \end{pmatrix}$$

X...can be set to zero, because only on surviving component can be actively projected



Generalized Bell inequality for kaons

$$S_{CHSH}(\bar{K}^0, \bar{K}^0, \bar{K}^0, \bar{K}^0; t_a, t_b, t_c, t_d) \\ = \left| E(\bar{K}^0, t_a; \bar{K}^0, t_b) - E(\bar{K}^0, t_a; \bar{K}^0, t_c) \right| + \left| E(\bar{K}^0, t_d; \bar{K}^0, t_b) + E(\bar{K}^0, t_d; \bar{K}^0, t_c) \right| \leq 2$$

II: vary in times

$$|k_n\rangle = \alpha |K^0\rangle + \beta |\bar{K}^0\rangle \dots \dots \text{quasi-spin}$$

"Are you an antikaon or not at a certain time?"

$$S_{Photon} = 2\sqrt{2}$$

$$S_{Kaon} \left(\frac{3\pi}{4}, \frac{3\pi}{4}, \frac{\pi}{2}, \mathbf{0} \right) = 1.36$$

Violation!

NO violation!

Strangeness oscillation/decay:

$$x = \frac{\Delta m}{\Gamma} \approx \frac{2 \Delta m}{\Gamma_s} \approx 1$$

PROPOSITION:

The CHSH-inequality is violated iff $x > 2$ for kaons.

Bell inequality sensitive to strangeness violated?

$$S_{CHSH}(\bar{K}^0, \bar{K}^0, \bar{K}^0, \bar{K}^0; t_a, t_b, t_c, t_d) \\ = \left| E(\bar{K}^0, t_a; \bar{K}^0, t_b) - E(\bar{K}^0, t_a; \bar{K}^0, t_c) \right| + \left| E(\bar{K}^0, t_d; \bar{K}^0, t_b) + E(\bar{K}^0, t_d; \bar{K}^0, t_c) \right| \leq 2$$

Can we violate the BI for a certain initial state and what is the maximum value?

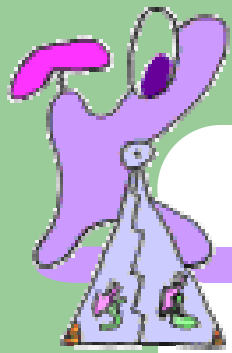
Arbitrary initial state:

$$|\psi\rangle = r_1 e^{i\varphi_1} |K_S K_S\rangle + r_2 e^{i\varphi_2} |K_S K_L\rangle + r_3 e^{i\varphi_3} |K_L K_S\rangle + r_4 e^{i\varphi_4} |K_L K_L\rangle$$

Expectation value:

$$E_{\bar{K}^0, \bar{K}^0}(t_l, t_r) = 1 + r_1^2 e^{-\Gamma_S(t_l+t_r)} + r_2^2 e^{-\Gamma_S t_l - \Gamma_L t_r} + r_3^2 e^{-\Gamma_L t_l - \Gamma_S t_r} + r_4^2 e^{-\Gamma_L(t_l+t_r)} \\ - r_1^2 (e^{-\Gamma_S t_l} + e^{-\Gamma_S t_r}) - r_2^2 (e^{-\Gamma_S t_l} + e^{-\Gamma_L t_r}) - r_3^2 (e^{-\Gamma_L t_l} + e^{-\Gamma_S t_r}) - r_4^2 (e^{-\Gamma_L t_l} + e^{-\Gamma_L t_r}) \\ + 2r_1 r_2 (1 - e^{-\Gamma_S t_l}) \cos(\Delta m t_r + \phi_1 - \phi_2) e^{-\Gamma t_r} + 2r_1 r_3 \cos(\Delta m t_l + \phi_1 - \phi_3) e^{-\Gamma t_l} (1 - e^{-\Gamma_S t_r}) \\ + 2r_2 r_4 \cos(\Delta m t_l + \phi_2 - \phi_4) e^{-\Gamma t_l} (1 - e^{-\Gamma_L t_r}) + 2r_3 r_4 (1 - e^{-\Gamma_L t_l}) \cos(\Delta m t_r + \phi_3 - \phi_4) e^{-\Gamma t_r} \\ + 2r_1 r_4 \cos(\Delta m(t_l + t_r) + \phi_1 - \phi_4) e^{-\Gamma(t_l+t_r)} + 2r_2 r_3 \cos(\Delta m(t_l - t_r) + \phi_2 - \phi_3) e^{-\Gamma(t_l+t_r)} .$$

$$E^{photons}(\vec{a}, \vec{b}) = 2r_1 r_4 \sin 2(\varphi_a - \varphi_b) + 2r_2 r_3 \cos 2(\varphi_a - \varphi_b)$$



Bell inequality sensitive to strangeness violated?

Can we violate the BI for a certain initial state and what is the maximum value?

Hiesmayr, European Journal C (2007) or
quant-ph/0607210

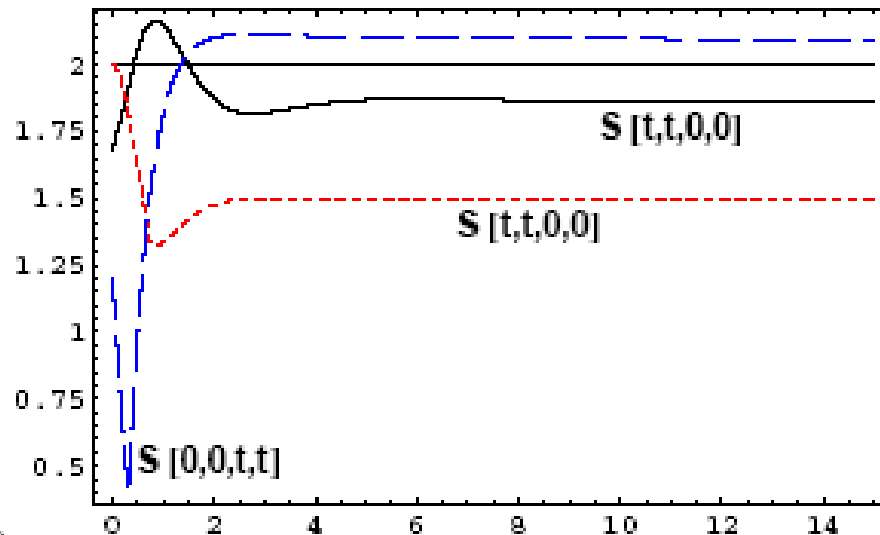
!! YES !!

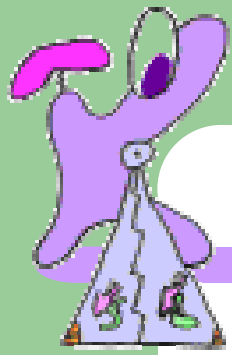
S=2.1596

χ : ($r_1=-0.7823$;
 $r_2=r_3=0.1460$; $r_4=0.5$
877; $\Phi_1=;$...;
 $t_1=t_2=1.79\tau_S$;
 $t_3=t_4=0$)

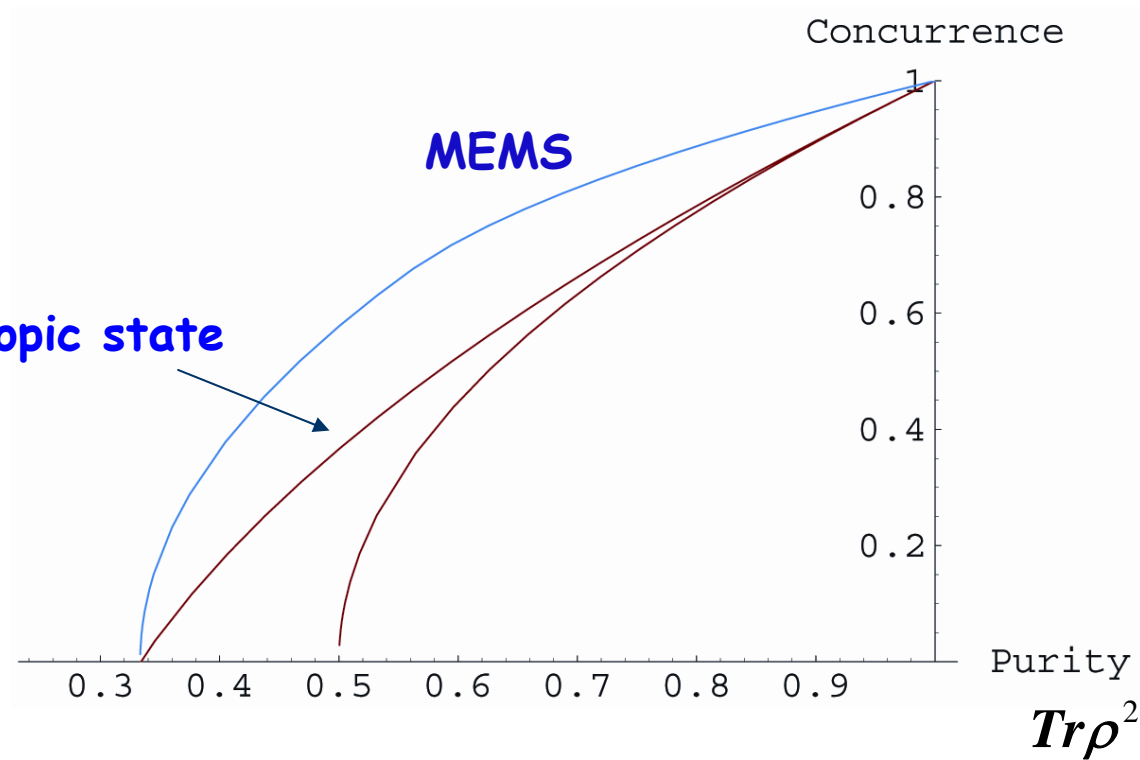
S=2.1175

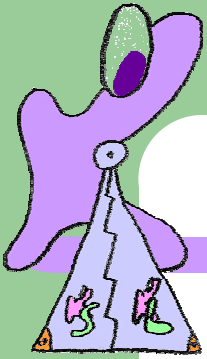
ξ : ($r_1=-0.8335$;
 $r_2=r_3=-0.2446$;
 $r_4=0.4415$; $t_1=t_2=0$;
 $t_3=t_4=5.77\tau_S$)





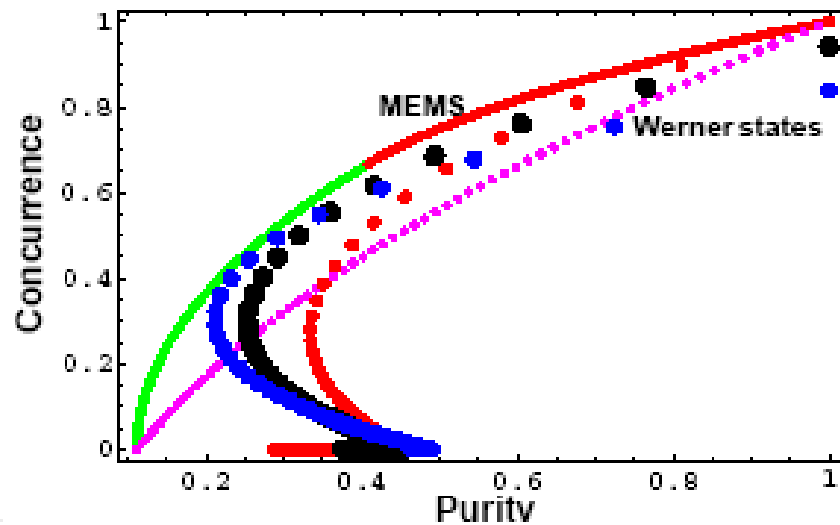
Maximally entangled maximally mixed states MEMS



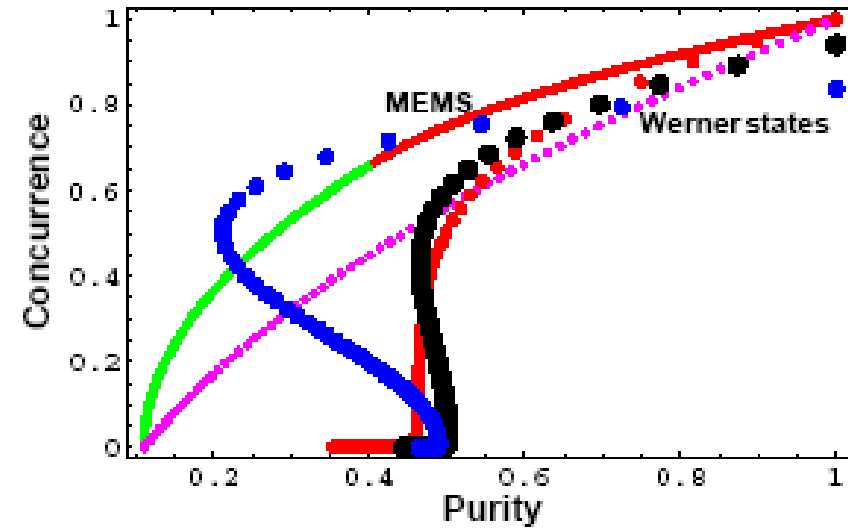


Entanglement—purity relation

$t_l = t_r = t$



$t_l = 0; t_r = t$



Normalized purity:
$$\frac{d \text{Tr} \rho^2 - 1}{d - 1}$$

Entanglement of formation = Concurrence for kaons

The generalized BI for kaons

Expectation value:

valid for any LRT and QM

$$|k_n\rangle = \alpha |K^0\rangle + \beta |\bar{K}^0\rangle \dots \dots \text{quasi-spin}$$

$$\begin{aligned} E(k_n, t_a; k_m, t_b) &= P_{n.m}(Y, t_a; Y, t_b) + P_{n.m}(N, t_a; N, t_b) \\ &\quad - P_{n.m}(Y, t_a; N, t_b) - P_{n.m}(N, t_a; Y, t_b) \\ &= -1 + 2\{P_{n.m}(Y, t_a; Y, t_b) + P_{n.m}(N, t_a; N, t_b)\} \end{aligned}$$

Generalized CHSH-Bell inequality

$$S(k_n, k_m, k_{n'}, k_{m'}; t_a, t_b, t_c, t_d) =$$

local realistic theories

$$|E(k_n, t_a; k_m, t_b) - E(k_n, t_a; k_{m'}, t_c)| + |E(k_{n'}, t_d; k_m, t_b) + E(k_{n'}, t_d; k_{m'}, t_c)| \leq 2$$

inequality
sensitive to
strangeness

S

inequality
sensitive to
CP violation in mixing

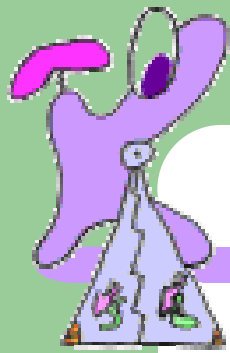
δ

inequality
sensitive to
CP violation in the
decay amplitudes

ϵ'

inequality
sensitive to
regeneration

$E[\tau, \delta\tau]$



Summary



- decay is a `kind of decoherence`
- Bell inequality sensitive to strangeness can be maximally violated for a non-maximally entangled state!

Nonlocality and entanglement are different quantum features!

Literature:

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