#### Coherent States in Physics and Mathematics - III

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Expository Quantum Lecture Series 5

Institute for Mathematical Research Putra University, Malaysia

Jan 9 - 13, 2012

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#### Abstract

In this lecture we look at some mathematical structures which will lead us to the definition of coherent states, in the general setting that we shall adopt. The definition will be broad enough to encompass all the different types of CS appearing in the physical and mathematical literature. Further generalizations, while possible and sometimes useful, will not be considered here.

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2 Some related mathematical structures



Image: A matrix and a matrix

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The canonical CS, that we just studied, were essentially a family of vectors in a Hilbert space, which satisfied a resolution of the identity and defined a POV measure. This fact turns out to be the mathematical property, crucial to any generalization of the concept. We start out, therefore, with a fairly rigorous treatment of positive operator valued measures and frames and reproducing kernels in that setting.

As before, let  $\mathfrak{H}$  be a separable, complex Hilbert space,  $\mathcal{L}(\mathfrak{H})$  the set of all bounded, linear operators on  $\mathfrak{H}$  and  $\mathcal{L}(\mathfrak{H})^+$  its subset of positive elements. Note that  $A \in \mathcal{L}(\mathfrak{H})^+$  iff A is self-adjoint and  $\langle \phi | A \phi \rangle \ge 0$ ,  $\forall \phi \in \mathfrak{H}$ . In particular, A = P is an (orthogonal) projection operator iff  $P = P^* = P^2$ .

Let X be a metrizable, locally compact space. (All the group spaces and parameter spaces for defining CS will be of this type. Metrizability is a technical assumption, which entails that Baire subsets and Borel subsets of X coincide). The Borel sets  $\mathcal{B}(X)$  of X consist of elements of the  $\sigma$ -algebra formed by its open sets.

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A positive Borel measure  $\nu$  on  $\mathcal{B}(X)$  is a map  $\nu : \mathcal{B}(X) \to \overline{\mathbb{R}^+} = \mathbb{R}^+ \cup \{\infty\}$  satisfying:

• For the empty set,  $\emptyset$ ,

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$$u(\emptyset) = 0.$$

If J is a countable index set and ∆<sub>i</sub>, i ∈ J, are mutually disjoint elements of B(X) (i.e., ∆<sub>i</sub> ∩ ∆<sub>j</sub> = Ø, for i ≠ j), then

$$u(\cup_{i\in J}\Delta_i) = \sum_{i\in J} 
u(\Delta_i)$$

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If in addition,  $\nu$  satisfies

$$u(\Delta) = \sup_{C \subseteq \Delta, \ C = \ compact} \nu(C),$$

then it is called a regular Borel measure. Unless the contrary is stated, all Borel measures will be assumed to be regular. The measure  $\nu$  is said to be bounded or finite if  $\nu(X) < \infty$ .

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Let  $\mathfrak{H}$  be a Hilbert space and X a locally compact space. A positive operator-valued (POV) measure is a map  $a : \mathcal{B}(X) \to \mathcal{L}(\mathfrak{H})^+$ , satisfying:

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The above sum is understood to converge weakly.

This means that, if  $\phi$  and  $\psi$  are arbitrary elements of  $\mathfrak{H}$ , then the complex sum  $\sum_{i \in I} \langle \phi | \mathbf{a}(\Delta_i) \psi \rangle$  converges.

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The POV-measure *a* is said to be regular if, for each  $\phi \in \mathfrak{H}$ , the measure  $\mu_{\phi}$  is regular. It is said to be bounded if

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This just means that each one of the positive Borel measures  $\mu_{\phi}$  above is bounded. In particular, the POV-measure is said to be normalized if

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Let  $F: X \to \mathcal{L}(\mathfrak{H})^+$  be a weakly measurable function, i.e., for each  $x \in X$ , F(x) is a bounded, positive operator, and for arbitrary  $\phi, \psi \in \mathfrak{H}$ , the complex function  $x \mapsto \langle \phi | F(x)\psi \rangle$  is measurable in  $\mathcal{B}(X)$ .

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Using F and the regular Borel measure  $\nu$ , we may define the regular POV-measure a, such that for all  $\Delta \in \mathcal{B}(X)$ ,

$$a(\Delta) = \int_{\Delta} F(x) \ d\nu(x)$$

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The operator function F is then called a density for the POV-measure a. We shall assume in all such cases that the support of the measure  $\nu$  is all of X. (This is not a severe restriction on the class of measures used, since one can always delete the complement of the support of  $\nu$  in X, and work with a smaller space X' which is just the support of  $\nu$ ).

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Furthermore, we shall mainly be concerned with bounded POV-measures with densities, such that the operator

$$A = \int_X F(x) \ d\nu(x)$$

has an inverse which is densely defined on  $\mathfrak{H}$ . (This ensures that using F it is possible to define a family of vectors which span  $\mathfrak{H}$ ).

The operator A is sometimes referred to as the resolution operator. We will not, in general, assume  $A^{-1}$  to be bounded. However, the following special case will be of particular interest to us.

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Assume that for each  $x \in X$ , the operators F(x) have the same finite rank n. Then there exists, for each x, a set of linearly independent vectors,  $\eta_x^i$ , i = 1, 2, ..., n, in  $\mathfrak{H}$  for which the map  $x \mapsto \eta_x^i$  is measurable and

$$F(x) = \sum_{i=1}^{n} |\eta_x^i\rangle\langle\eta_x^i|$$

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A set of vectors  $\eta_x^i$ , i = 1, 2, ..., n,  $x \in X$ , constitutes a rank-*n* frame, denoted  $\mathcal{F}{\eta_x^i, A, n}$ , if for each  $x \in X$ , the vectors  $\eta_x^i$ , i = 1, 2, ..., n, are linearly independent and if they satisfy the above integral relation, with both *A* and  $A^{-1}$  being bounded positive operators on  $\mathfrak{H}$ .

If  $A = \lambda I$ , for some  $\lambda > 0$ , the frame is called tight and exact if  $\lambda = 1$ .

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The frame condition may alternatively be written in the following more familiar form:

$$\mathsf{m}(A) \|f\|^2 \leq \sum_{i=1}^n \int_X |\langle \eta_x^i | f \rangle|^2 \, d\nu(x) \leq \mathsf{M}(A) \|f\|^2, \, \forall f \in \mathfrak{H}$$

where m(A) and M(A) denote, respectively, the infimum and the supremum of the spectrum of A.

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These numbers, which satisfy  $0 < m(A) \le M(A) < \infty$ , are usually called the frame bounds.

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All this takes a more familiar shape if the space X is discrete, with  $\nu$  a counting measure. Indeed the frame relation relation then reads:

$$A = \sum_{i=1}^{n} \sum_{x \in X} |\eta_x^i\rangle \langle \eta_x^i|,$$

and the frame is discrete.

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#### Theorem (Naimark's extension theorem)

Let a be an arbitrary normalized POV-measure on  $\mathcal{B}(X)$ . Then there exist a Hilbert space  $\tilde{\mathfrak{H}}$ , an isometric embedding  $W : \mathfrak{H} \to \tilde{\mathfrak{H}}$  and a PV-measure  $\tilde{P}$  on  $\mathcal{B}(X)$ , with values in  $\mathcal{L}(\tilde{\mathfrak{H}})$ , such that  $\forall \Delta \in \mathcal{B}(X)$ ,

 $W_a(\Delta)W^{-1} = \mathbb{P}\widetilde{P}(\Delta)\mathbb{P}, \quad \mathbb{P} = projection \ operator: \quad \mathbb{P}\widetilde{\mathfrak{H}} = W\mathfrak{H},$ 

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(with  $W^{-1}$  defined on the range of W). The extension  $\{\widetilde{P}, \widetilde{\mathfrak{H}}\}$  of  $\{a, \mathfrak{H}\}$  can be chosen to be minimal in the sense that the following set of vectors be dense in  $\widetilde{\mathfrak{H}}$ :

 $\mathcal{S} = \{\widetilde{P}(\Delta) \Phi \mid \Delta \in \mathcal{B}(X), \Phi \in W\mathfrak{H}\}$ 

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This minimal extension is unique up to unitary equivalence.

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### The case of an exact frame

For illustrative purposes, we work out the Naimark extension for the cases where the POV measure *a* defines an exact frame, (i.e., A = I).

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$$\|\Phi\|^2 = \sum_{i=1}^n \int_X |\Phi^i(x)|^2 \ d\nu(x) < \infty,$$

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 $\Phi^i(x)$  denoting the *i*-th component of  $\Phi(x)$  in  $\mathbb{C}^n$ . Let  $W : \mathfrak{H} \to \widetilde{\mathfrak{H}}$  be the linear map,

$$(W\phi)^i(x) = \Phi^i(x) = \langle \eta^i_x | \phi \rangle, \quad i = 1, 2, \dots, n, \quad x \in X.$$

Then W is an isometry. Indeed,

$$\sum_{i=1}^n \int_X |\Phi^i(x)|^2 \ d\nu(x) = \sum_{i=1}^n \int_X \langle \phi | \eta^i_x \rangle_{\mathfrak{H}} \langle \eta^i_x | \phi \rangle_{\mathfrak{H}} \ d\nu(x),$$

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which, in view of the weak convergence of the integral, implies that

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If  $\Phi$  is an arbitrary vector in the range,  $W\mathfrak{H}$ , of  $\mathfrak{H}$  in  $\widetilde{\mathfrak{H}}$ , it is easily verified that the inverse map  $W^{-1}$  acts on it as

$$W^{-1}\Phi = \sum_{i=1}^n \int_X \Phi^i(x) \eta^i_x \, d\nu(x) \in \mathfrak{H},$$

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giving a reconstruction formula for  $\phi = W^{-1}\Phi$ , when the coefficients  $\Phi^i(x) = \langle \eta^i_x | \phi \rangle$ , i = 1, 2, ..., n,  $x \in X$  are known.

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giving a reconstruction formula for  $\phi = W^{-1}\Phi$ , when the coefficients  $\Phi^i(x) = \langle \eta^i_x | \phi \rangle$ , i = 1, 2, ..., n,  $x \in X$  are known. On  $\tilde{\mathfrak{H}}$  define the projection operators  $\tilde{P}(\Delta), \Delta \in \mathcal{B}(X)$ , by

$$(\widetilde{P}(\Delta)\Phi)(x) = \chi_{\Delta}(x)\Phi(x), \quad \Phi \in \widetilde{\mathfrak{H}},$$

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where  $\chi_{\Delta}$  is the characteristic function of the set  $\Delta$ . Clearly  $\widetilde{P}$  is a PV-measure on  $\mathcal{B}(X)$  with values in  $\mathcal{L}(\widetilde{\mathfrak{H}})$ .

We then have the result:

Theorem

The PV-measure  $\tilde{P}$  extends the POV-measure,

$$a(\Delta) = \sum_{i=1}^n \int_\Delta |\eta^i_x\rangle \langle \eta^i_x| \ d
u(x), \qquad \Delta \in \mathcal{B}(X),$$

defining the exact frame  $\mathcal{F}\{\eta_x^i, I, n\}$ , minimally in the sense of Naimark.

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#### Vector coherent states and coherent states

We now state what we will adopt as the mathematical definition of a family of coherent states. Let  $\mathcal{F}\{\eta_x^i, I, n\}$  be a rank-*n* exact frame on a Hilbert space  $\mathfrak{H}$ , with  $x \in X$ , a metrizable, locally compact space, equipped with a regular Borel measure  $\nu$ . Then,

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with such a family there is always a normalized POV measure,

$$a(\Delta) = \int_{\Delta} F(x) d\nu(\Delta) , \quad \Delta \in \mathcal{B}(\mathfrak{H}) , \quad \text{with} \quad a(X) = I ,$$

admitting a density function,

$$F(x) = \sum_{i=1}^n |\eta_x^i\rangle\langle \eta_x^i| , \quad x \in X .$$

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From our previous discussion, we see that given a family of vector CS, there is a canonical isometry

 $W: \mathfrak{H} \mapsto \widetilde{\mathfrak{H}} = \mathbb{C}^n \otimes L^2(X, d\nu) , \quad (W\phi)^i(x) := \Phi^i(x) = \langle \eta^i_x \mid \phi \rangle ,$  $\forall x \in X , \quad i = 1, 2, \dots, n .$ 

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The range of this isometry,  $\mathfrak{H}_{\mathcal{K}} := W\mathfrak{H}$ , is a closed subspace of  $\mathfrak{H}$  and we shall use the notation,

$$\boldsymbol{\xi}_{x}^{i}(\boldsymbol{y}) := (W\eta_{x}^{i})(\boldsymbol{y}) = \begin{pmatrix} \boldsymbol{\xi}_{x}^{i1}(\boldsymbol{y}) \\ \boldsymbol{\xi}_{x}^{i2}(\boldsymbol{y}) \\ \vdots \\ \boldsymbol{\xi}_{x}^{in}(\boldsymbol{y}) \end{pmatrix} \in \mathbb{C}^{n}, \quad \boldsymbol{y} \in \boldsymbol{X}, \quad i = 1, 2, \dots, n$$

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$$\xi_x^{ij}(y) = \langle \boldsymbol{\xi}_y^j \mid \boldsymbol{\xi}_x^i \rangle , | \quad i,j = 1, 2, \dots, n .$$

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In other words, as the name implies,  $y \mapsto \boldsymbol{\xi}_x^i(y)$ ,  $y \in X$ , is an *n*-vector valued function on X, which are well-defined for all  $y \in X$ . The reason for the subscript K in  $\mathfrak{H}_K$  will soon be clarified.

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If we denote by  $\mathbb{P}_{\mathcal{K}}$  the projection operator,  $\mathbb{P}_{\mathcal{K}}\widetilde{\mathfrak{H}} = \mathfrak{H}_{\mathcal{K}}$ , then in view of the fact that W is an isometry, we have the resolution of the identity on  $\mathfrak{H}_{\mathcal{K}}$ 

$$\sum_{i=1}^n \int_X |m{\xi}^i_x
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and the corresponding normalized POV measure on  $\mathfrak{H}_{\mathcal{K}}$ ,

$$\mathsf{a}_{\mathsf{K}}(\Delta):=\mathsf{W}\!\mathsf{a}(\Delta)\mathsf{W}^{-1}=\sum_{i=1}^n\int_\Delta|\boldsymbol{\xi}_x^i
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Thus, a family of vector CS can always be written as a set of everywhere defined, bounded vector-valued functions. Boundedness follows from

$$|\xi_x^{ij}(y)| = |\langle \boldsymbol{\xi}_y^j \mid \boldsymbol{\xi}_x^i \rangle| \le \|\boldsymbol{\xi}_y^j\| \|\boldsymbol{\xi}_x^i\|.$$

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The Hilbert space

$$\mathfrak{H}_{\mathcal{K}} = W\mathfrak{H} = \mathbb{P}_{\mathcal{K}}\widetilde{\mathfrak{H}} = \mathbb{P}_{\mathcal{K}}[\mathbb{C}^n \otimes L^2(X, d\nu)]$$

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Indeed, consider the  $n \times n$ -matrix valued function **K** on  $X \times X$ , with matrix elements

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This function defines a positive definite kernel, satisfying

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$$\int_{X} \mathbf{K}(x, z) \mathbf{K}(z, y) \, d\nu(z) = \mathbf{K}(x, y)$$

(matrix multiplication being implied in the integrand.)

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This last condition reflects the fact that  $\mathbb{P}^2_{\mathcal{K}} = \mathbb{P}_{\mathcal{K}}$ , and that  $\mathcal{K}(x, y)$  is the integral kernel of  $\mathbb{P}_{\mathcal{K}}$ . Indeed,  $\forall \Phi \in \widetilde{\mathfrak{H}} = \mathbb{C}^n \otimes L^2(X, d\nu)$  we have,

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Note that  $\Phi(y)$  is an *n*-vector valued function. This relation is easily verified to be a consequence of the resolution of the identity condition satisfied by the vector coherent states  $\boldsymbol{\xi}_x^i$  (which thus form an overcomplete set in  $\mathfrak{H}_K$ ).

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Moreover, the above relation also implies that if  $\Phi\in\mathfrak{H}_{\mathcal{K}},$  then

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It ought to be stressed that the above relation holds for all  $x \in X$  and not just up to a set of measure zero. Indeed, the functions in  $\mathfrak{H}_K$ , as stressed earlier, are well defined for all  $x \in X$ .

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$$oldsymbol{\xi}^i_{ imes} = \sum_{n=1}^d c^i_n(x) oldsymbol{\Psi}_n \ , \qquad c^i_n(x) \in \mathbb{C} \ .$$

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Alternatively, we may write,

$$\mathbf{K}(x,y) = \sum_{n=1}^{d} \mathbf{\Psi}_{n}(x) \mathbf{\Psi}_{n}(y)^{\dagger} .$$

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Furthermore, it is easy to see from the above equation that the inequality

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It also follows from this inequality that the elements of  $\mathfrak{H}_{\mathcal{K}}$  are functions which are everywhere well defined.

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implying that

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implying that

$$\begin{split} \|\Phi(x)\|_{\mathbb{C}^{n}}^{2} &= \sum_{i=1}^{N} |\Phi^{i}(x)|^{2} \leq \left[\sum_{i=1}^{N} \|\eta_{x}^{i}\|_{\mathfrak{H}}^{2}\right] \|\phi\|_{\mathfrak{H}}^{2} \\ \mathsf{But} \ \|\eta_{x}^{i}\|_{\mathfrak{H}}^{2} &= \|\boldsymbol{\xi}_{x}^{i}\|_{\mathfrak{H}^{K}}^{2} = \mathcal{K}_{ii}(x,x) \text{ and } \|\phi\|_{\mathfrak{H}}^{2} = \|\Phi\|_{\mathfrak{H}^{K}}^{2}, \text{ so that} \\ \|\Phi(x)\|_{\mathbb{C}^{n}} &\leq \left[\sum_{i=1}^{N} \mathcal{K}_{ii}(x,x)\right]^{\frac{1}{2}} \|\Phi\|_{\mathfrak{H}_{K}} = \mathsf{Tr}\left[\mathsf{K}(x,x)\right]^{\frac{1}{2}} \|\Phi\|_{\mathfrak{H}_{K}} \end{split}$$

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#### This is the statement that for each $x \in X$ , the evaluation map

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Thus, the reproducing kernel Hilbert space  $\mathfrak{H}_{\mathcal{K}}$  is a space in which the evaluation map is continuous at each point.

Indeed, continuity of the evaluation map guarantees the existence of coherent states and a reproducing kernel.
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- 4. Functions in this reproducing kernel Hilbert space are pointwise well-defined and the corresponding evaluation map is continuuous.
- 5. The vectors  $\boldsymbol{\xi}_x^i$  or  $\eta_x^i$ , as defined here, are not necessarily normalized. However, physical coherent states will always be normalized, i.e., physically, we shall work with the vectors  $\boldsymbol{\xi}_x^i/\|\boldsymbol{\xi}_x^i\|$  or  $\eta_x^i/\|\eta_x^i\|$ .

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