

Arithmetic Quantum Chaos

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Expository Quantum Lecture Series 5
Universiti Putra Malaysia

January 12, 2012

Part I

Introduction and Motivation

A Classical Problem



- Free point particle on a Riemann surface (orbifold) \mathcal{M} .
- Geodesic flow: $\Phi : T^1\mathcal{M} \rightarrow T^1\mathcal{M}$. (C)
- $T^1\mathcal{M} \simeq \mathcal{M} \times S^1$, the unit tangent (or co-tangent) bundle of \mathcal{M} .

Quantization

- Quantization \Rightarrow The Schrödinger equation
- Separation of variables \Rightarrow
Eigenvalue problem for Δ – the Laplacian on \mathcal{M}

$$(\hbar^2 \Delta + \tilde{\lambda}_k \psi_k) = 0 \Leftrightarrow$$

$$\boxed{(\Delta + \lambda_k) \psi_k = 0} \tag{Q}$$

- We can choose an orthonormal basis $\{\psi_k\}$ and associated eigenvalues $\lambda_1 \leq \lambda_2 \leq \dots$ (finite or infinite)
- **Main Question:** Is it possible to tell from properties of the ψ_k or the λ_k whether the classical system is chaotic?

Integrable v/s non-Integrable

- If (\mathcal{M}, Φ_t) is integrable, then quantization is well-understood (Bohr-Sommerfeld, Einstein)
- If completely integrable then ψ_k localize on invariant torii and λ_k are uniformly distributed.
 - Explicit quantization of the form:

$$\int p dq = 2\pi\hbar l_k$$

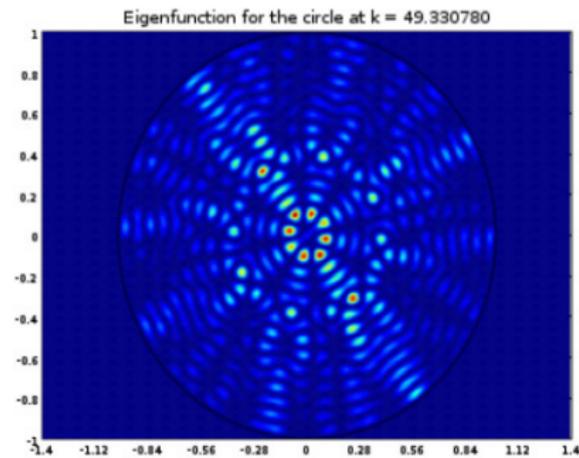
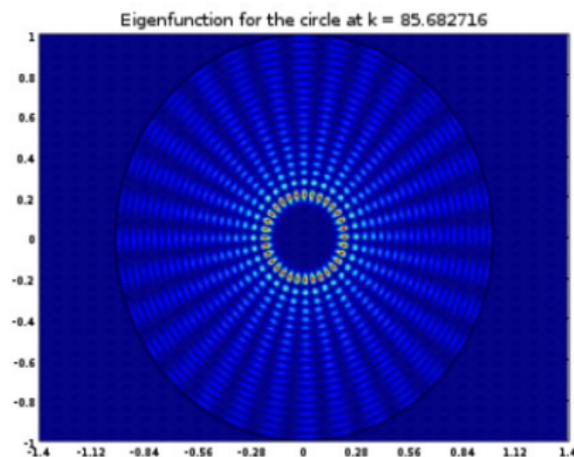
- For generic non-integrable systems it is not possible (or unknown how) to find such explicit quantization conditions.
- We are interested in chaotic (ergodic, Anosov etc.) systems.

Properties of the Quantum System

- Value-distribution
 - quantum ergodicity
 - Berry's random wave model (on small scales) etc.
- Level spacing
 - Random matrix theory

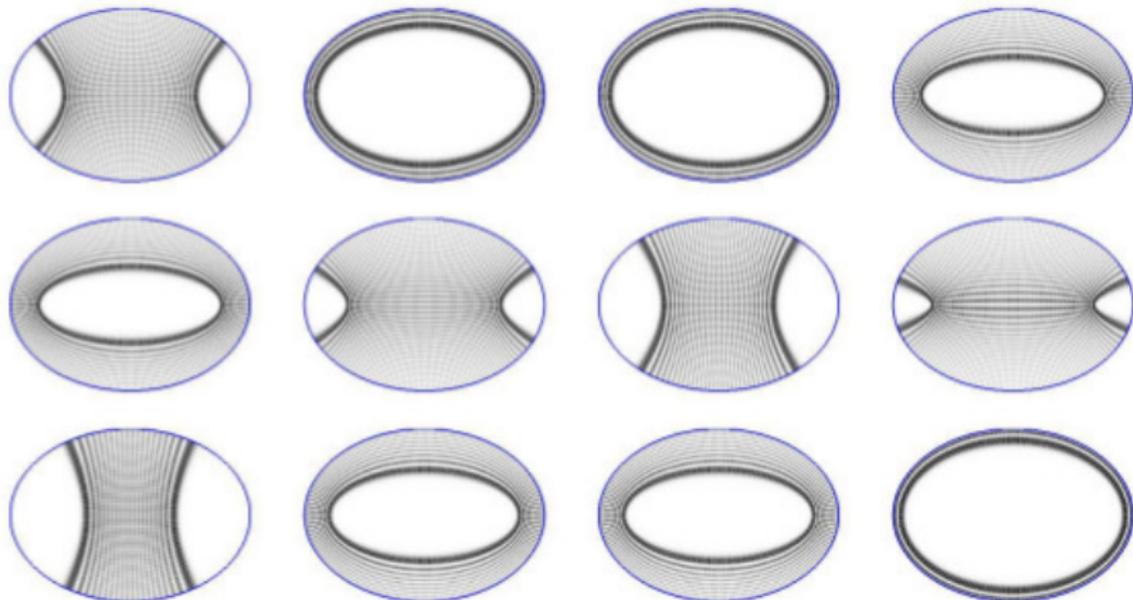
We will focus on the first property in this lecture.

Integrable Example



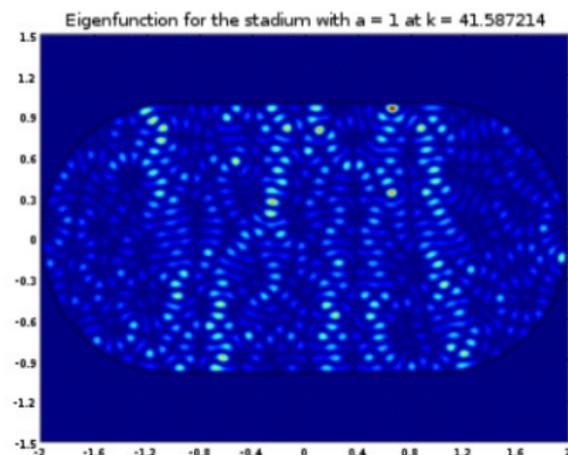
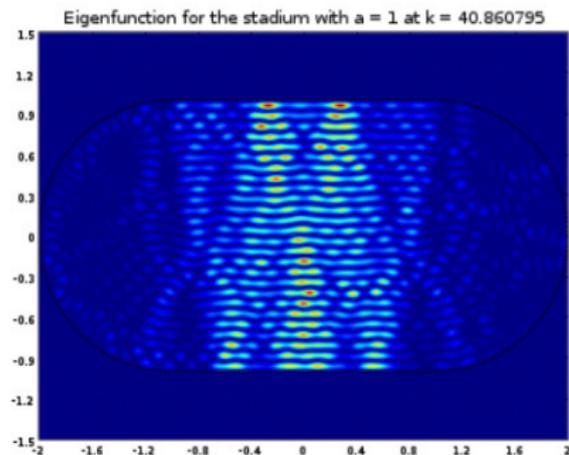
(Bober, „Eigenvalue Statistics for some Quantum Billiards“)

Integrable Example 2



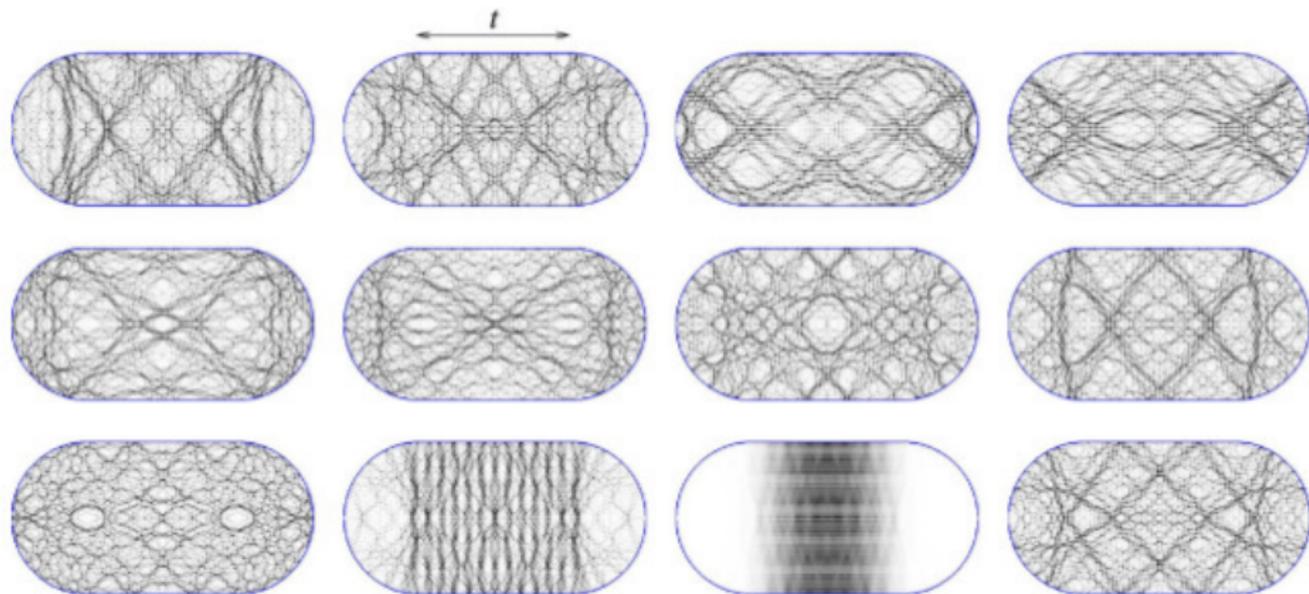
(Sarnak, „Recent Progress on QUE“, computed by Barnett)

Chaotic Example: Stadium



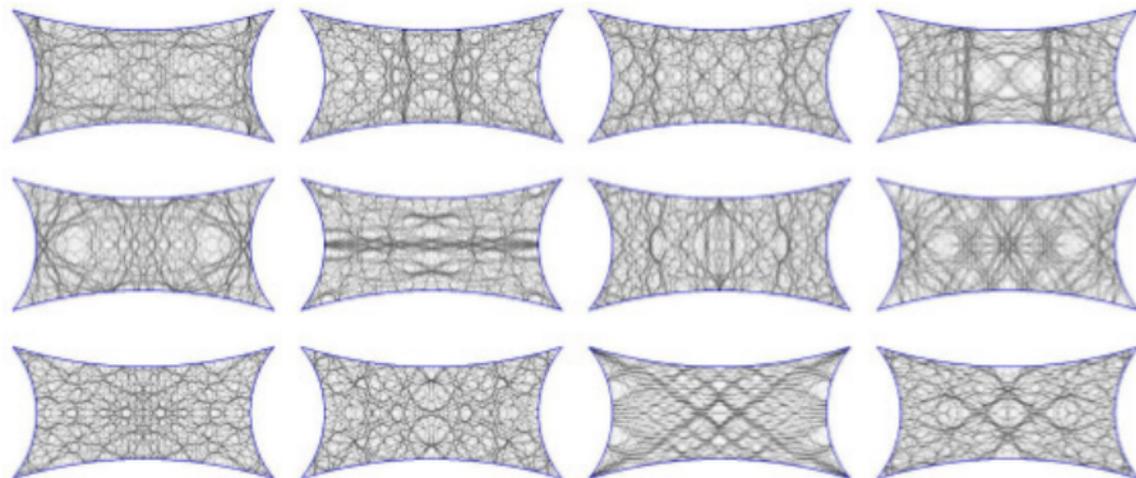
(Bober, „Eigenvalue Statistics for some Quantum Billiards“)

Chaotic Example: Stadium 2



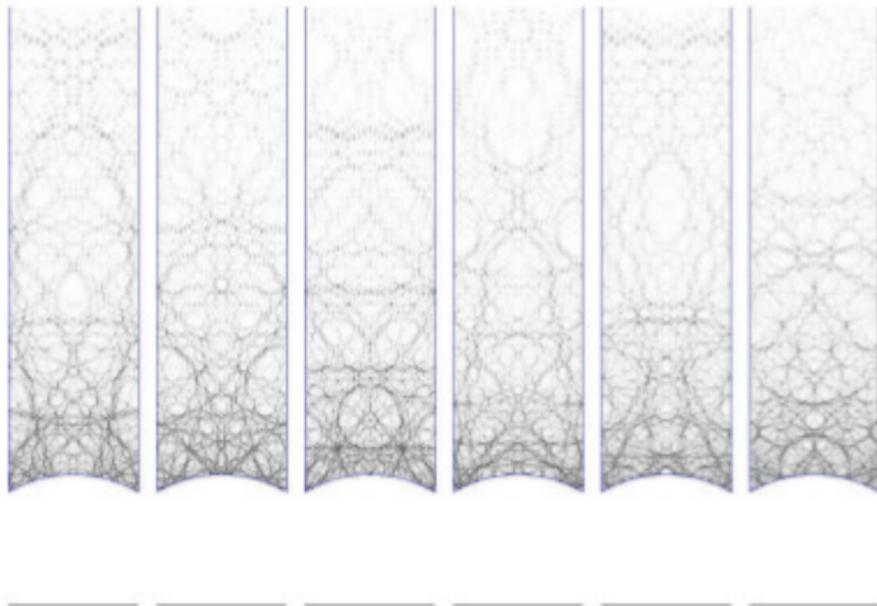
(Sarnak, „Recent Progress on QUE“, computed by Barnett)

Chaotic Example: Dispersing Sinai Billiard



(Sarnak, „Recent Progress on QUE“, computed by Barnett)

Chaotic Example: The modular surface



(From Sarnak, „Recent Progress on QUE“, computed by Then)

Quantum Ergodicity

- Let $\{\psi_k\}$ be an ON-basis of Δ -eigenfunctions of $L^2(\mathcal{M}, d\mu)$ ($d\mu$ is the area measure on \mathcal{M}) with associated eigenvalues $\lambda_1 \leq \lambda_2 \leq \dots$ (finite or infinite)
- Define a measure

$$\nu_\psi = |\psi(z)|^2 d\mu(z)$$

- Micro-local lift: μ_ψ – a measure on $T_1\mathcal{M}$

$$\mu_\psi(f) = \langle \text{Op}(f)\psi, \psi \rangle$$

- Check: If $f(z, \xi) := f(z)$ then $\text{Op}(f)\psi(z) = f(z)\psi(z) \Rightarrow$ restriction of $d\mu_\psi$ is ν_ψ .

Quantum Ergodicity

- Any weak limit of ψ_k is a probability measure, such a measure is called a *quantum limit*.
- Schnirelman: Any quantum limit must be invariant under the geodesic flow.
- If $\Phi_t : \mathcal{M} \rightarrow \mathcal{M}$ is ergodic then we know that:
 - μ -almost all geodesics are μ -equidistributed in $T^1\mathcal{M}$ (Birkhoff's theorem)
 - *Quantum Ergodicity (QE)* holds, i.e. almost all quantum limits are equal to the Liouville measure μ .

Remark

If all quantum limits are equal to μ we say that Quantum Unique Ergodicity (QUE) holds.

Possible Quantum Limits

- Which quantum limits can occur?
- Obvious candidates (invariant under geodesic flow):
 - The Liouville measure μ
 - Arc measure ds supported on a closed geodesic.
(Colin-de-Verdière and Zelditch)
- The second option is called „strong scars“.
- No scars on arithmetic surfaces (Rudnick and Sarnak, 1994)
- No scars on any compact negatively curved surface
(Anantharaman, 2008)

Quantum Unique Ergodicity (QUE)

Conjecture (Rudnick-Sarnak)

If \mathcal{M} is a negatively curved surface then Quantum Unique Ergodicity holds.

- The stadium is not QUE for almost all side-lengths (Hassell, 2008)
- QUE holds for quantum cat maps
 - along certain subsequences (Degli Eposti, Graffi, Isola, 1995)
 - for Hecke eigenbases (Kurlberg-Rudnick, 2001)

Arithmetic Quantum Unique Ergodicity

QUE holds for

- Hecke eigenstates on compact arithmetic surfaces
(Lindenstrauss, 2006)
- continuous spectrum for non-compact arithmetic surfaces
(Luo-Sarnak 1995 and Jakobson, 1994)
- discrete spectrum for non-compact arithmetic surfaces
(Soundararajan, 2009)

The main topic of these lectures will be to develop the necessary theory to understand these results.

Part II

Arithmetic Groups and Surfaces

Arithmetic Surfaces



- $\mathcal{M} \simeq \Gamma \backslash \mathcal{H}$, where
- \mathcal{H} is the hyperbolic upper half-plane.
- Γ is an arithmetic Fuchsian group.

Some Hyperbolic Geometry



- $\mathcal{H} = \{z = x + iy \in \mathbb{C} \mid y > 0\}$ with

$$ds = y^{-1} |dz|,$$

$$d\mu = y^{-2} dx dy.$$

- $\text{Isom}^+(\mathcal{H}) \simeq PSL_2(\mathbb{R}) \simeq SL(2, \mathbb{R}) / \{\pm 1\}$.

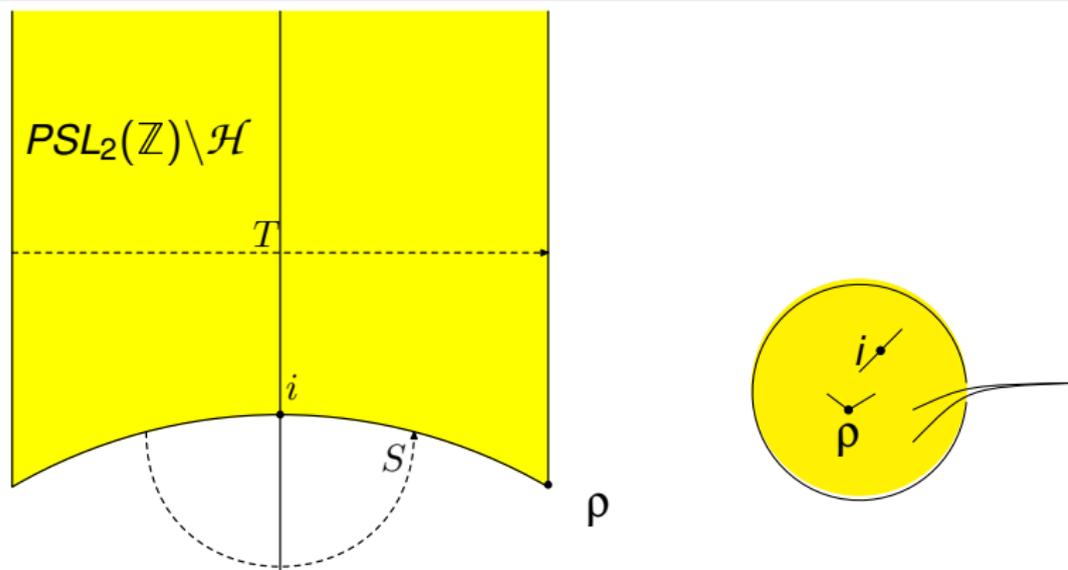
$$\gamma = \begin{pmatrix} a & b \\ c & d \end{pmatrix} : z \mapsto \frac{az + b}{cz + d}$$

- A *Fuchsian group* is a discrete subgroup of $PSL_2(\mathbb{R})$.

Arithmetic groups

- A Fuchsian group Γ is arithmetic if
 - $\Gamma = PSL_2(\mathbb{Z}) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix}, a, b, c, d \in \mathbb{Z}, ad - bc = 1 \right\}$
(non-compact), or
 - Γ is derived from a quaternion algebra (compact), or
 - Γ is commensurable with another arithmetic group.
- Important property: Infinitely many symmetries.
- Examples of important subgroups:
 - $\Gamma(N) = \{A \in PSL_2(\mathbb{Z}) \mid A \equiv 1_2 \pmod{N}\}$
(principal congruence subgroup)
 - $\Gamma_0(N) = \left\{ \begin{pmatrix} a & b \\ Nc & d \end{pmatrix}, a, b, c, d \in \mathbb{Z}, ad - bNc = 1 \right\}$

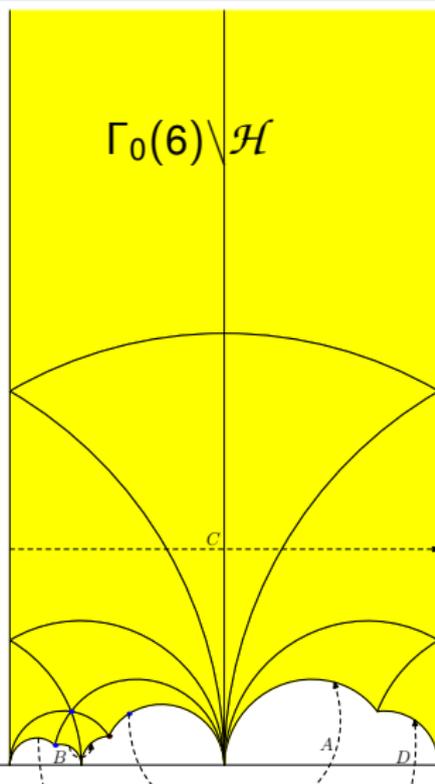
Fundamental domain for $PSL_2(\mathbb{Z})$



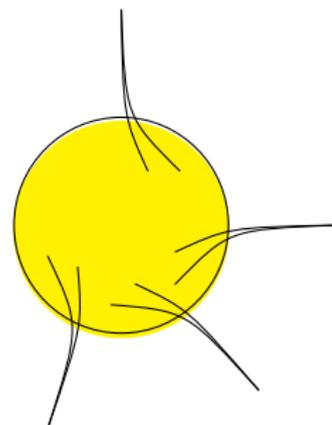
$$S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, z \mapsto -\frac{1}{z}$$

$$T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, z \mapsto z + 1$$

Fundamental domain for $\Gamma_0(6)$



→
identifying sides



The modular surface

We study functions on $X = \Gamma \backslash \mathcal{H}$ where $\Gamma = PSL_2(\mathbb{Z})$.

- $\text{Aut}(\Gamma)$ consists of $\varphi : \mathcal{H} \rightarrow \mathcal{H}$ satisfying

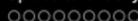
$$\varphi(\gamma z) = \varphi(z), \quad \forall \gamma \in \Gamma$$

- Petersson inner-product (φ, ψ measurable)

$$\langle \varphi, \psi \rangle = \int_X \varphi(z) \overline{\psi(z)} d\mu(z)$$

- $L^2(\Gamma)$ consists of $\varphi : \mathcal{H} \rightarrow \mathbb{C}$ measurable and

$$\| \varphi \|_2 = \langle \varphi, \varphi \rangle < \infty$$



Hyperbolic Laplace-Beltrami operator

Hyperbolic Laplacian

$$\Delta = y^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) = 4y^2 \frac{\partial^2}{\partial z \partial \bar{z}}.$$

- Δ is invariant under $PSL_2(\mathbb{R})$
- Δ can be extended to a self-adjoint non-positive elliptic differential operator on a dense subspace of $L^2(\Gamma)$.
- If $(\Delta + \lambda)\varphi = 0$ with $\varphi \in L^2(\Gamma)$ then $\lambda \geq 0$ and we write

$$\lambda = s(1 - s)$$

with $s \in \frac{1}{2} + i\mathbb{R}$ or $s \in [0, 1]$.

Automorphic Eigenfunctions

- Any $\varphi \in \text{Aut}(\Gamma)$ can be written

$$\varphi(z) = \sum_n c_n(y) e(nx), \quad e(x) = e^{2\pi i x}.$$

- If $(\Delta + \lambda)\varphi = 0$ then for $n \neq 0$:

$$c_n(y) = \alpha \sqrt{y} I_{s-\frac{1}{2}}(2\pi |n| y) + \beta \sqrt{y} K_{s-\frac{1}{2}}(2\pi |n| y)$$

(modified Bessel functions of the first and second kind)

- $K_s(y) = \int_0^\infty e^{-x \cosh t} \cosh(st) dt \rightarrow 0$ as $y \rightarrow \infty$,
- $I_s(y) = \sum_0^\infty \frac{1}{m! \Gamma(m+\alpha+1)} \left(\frac{x}{2}\right)^{2m+\alpha} \rightarrow \infty$ as $y \rightarrow \infty$.

If $n = 0$ then

$$c_0(y) = \begin{cases} ay^s + by^{1-s}, & s \neq \frac{1}{2}, \\ a\sqrt{y} + b\sqrt{y} \ln y, & s = \frac{1}{2}. \end{cases}$$



Automorphic eigenfunction

- Polynomial growth \Rightarrow

$$\varphi(z) = ay^s + by^{1-s} + \sum_{n \neq 0} c_n \sqrt{y} K_{s-\frac{1}{2}}(2\pi|n|y) e(nx).$$

- $L^2 \Rightarrow s \in \frac{1}{2} + i\mathbb{R} \cup [0, 1]$ and

- $s \in [0, \frac{1}{2}] \Rightarrow b = 0,$
- $s \in [\frac{1}{2}, 1] \Rightarrow a = 0$
- $s \in \frac{1}{2} + i\mathbb{R} \Rightarrow a = b = 0$

For $PSL_2(\mathbb{Z})$ $s \in \frac{1}{2} + i\mathbb{R}.$

Discrete Spectrum

Let $\{\psi_k\}$ and $\{\lambda_k\}$ be the set of L^2 -eigenfunctions. Then

- λ_j are discrete ,
- ψ_j are orthogonal (after diagonalizing any multiple eigenspaces)

Two types of eigenvalues in $[\frac{1}{2}, 1]$ (notice that $s \rightarrow 1 - s$ leaves λ invariant):

- Cuspidal eigenvalues (*exceptional*)
- Residual eigenvalues

None of these exist for the modular group: $\lambda_1 \geq \frac{3}{2}\pi^2$
(elementary estimates)

Spectral theory of the modular group

- $PSL_2(\mathbb{Z}) \backslash \mathcal{H}$ is non-compact \Rightarrow discrete and continuous spectrum
- Discrete spectrum: *Maass waveforms*, $\mathcal{M}(\Gamma)$
 - $\varphi(\gamma z) = \varphi(z)$, $\forall \gamma \in PSL_2(\mathbb{Z})$,
 - $(\Delta + \lambda)\varphi(z)$ for some $\lambda \geq 0$,
 - $\int_X |\varphi(z)|^2 d\mu(z) < \infty$
- We have

$$\varphi(z) = \sum_{n \neq 0} c_n \sqrt{y} K_{iR}(2\pi |n| y) e(nx)$$

$\varphi(z) \rightarrow 0$ as $y \rightarrow \infty$, i.e. φ is *cuspidal*.

The Continuous Spectrum

An elementary eigenfunction is y^s and we can form the *Eisenstein series*:

$$E(z; s) = \sum_{\gamma \in \Gamma_\infty \backslash \Gamma} (\Im \gamma z)^s, \quad \Re s > 1$$

where $\Gamma_\infty = \langle T \rangle = \{z \mapsto z + k \mid k \in \mathbb{Z}\}$ is the stabilizer of ∞ .

- $E(z; s)$ has an analytic continuation in z for each $\Re s > 1$.
- $E(z; s)$ has meromorphic continuation in s to \mathbb{C} without poles in $\Re s > \frac{1}{2}$.
- $E(z; s) = \varphi(s) E(z; 1 - s)$



The Eisenstein Series

- Fourier expansion

$$E(z; s) = y^s + \varphi(s) y^{1-s} + \sum \varphi_n(s) \sqrt{y} K_{s-\frac{1}{2}}(2\pi |n| y) e(nx).$$

with

$$\varphi(s) = \sqrt{\pi} \frac{\Gamma(s - \frac{1}{2}) \zeta(2s - 1)}{\Gamma(s) \zeta(2s)},$$

$$\varphi_n(s) = \frac{2\pi^s |n|^{s-\frac{1}{2}} \sigma_{1-2s}(|n|)}{\Gamma(s) \zeta(2s)}$$

- The function $\varphi(s)$ is called the scattering determinant
 - y^s is an incoming plane wave from $i\infty$,
 - y^{1-s} is the outgoing wave
- $\varphi(s) \varphi(1-s) = 1$.

Spectral decomposition of L^2

Theorem

If $f \in L^2(\Gamma)$ then

$$f(z) = \sum_k \langle f, \psi_k \rangle \psi_k(z) + \int_0^\infty g(t) E\left(z; \frac{1}{2} + it\right) dt$$

where

$$g(t) = \frac{1}{2\pi} \int_{\Gamma \setminus \mathcal{H}} f(\tau) \overline{E\left(\tau; \frac{1}{2} + it\right)} d\mu(\tau).$$

Level Spacings

Set $\lambda_j = \frac{1}{4} + R_j^2$

$$R_1 \leq R_2 \leq \dots$$

- Counting function

$$N(T) = \#\{j : R_j \leq T\}$$

- Weyl's law:

$$N(T) = \frac{\mu(\Gamma \backslash \mathcal{H})}{4\pi} T^2 - \kappa \frac{1}{\pi} T \ln T + O(T)$$

where κ is the number of cusps of Γ .

- Mean spacing

$$\delta \simeq \frac{1}{N'(T)} \sim \frac{2\pi}{\mu(\Gamma \backslash \mathcal{H}) T}.$$

Normalized spacing distribution

- Normalize eigenvalues by $t\gamma$

$$\tilde{R}_j = \frac{1}{\delta} R_j$$

Remark

In dimension 2: If there is a stable geodesic then there is an arithmetic progression of density T (with eigenfunns localizing along the geodesic)

- Consider the spacings $\delta_j = \tilde{R}_j - \tilde{R}_{j+1}$ as random numbers in $[0, \infty)$ with mean 1 and densityfn. $\rho(x)$

Conjectures

Conjecture

The eigenvalues of an integrable system follow a Poisson distribution.

(Wigner, Dyson, Mehta, Bohigas, Berry-Tabor etc.)

Conjecture

The eigenvalues of a quantized chaotic system behaves like eigenvalues of random matrices.

(Bohigas, Gianonni and Schmidt)

In both cases, we assume appropriate normalizations.

Standard Models

- Poisson: λ_j are given by a Poisson process

$$\rho_{\text{Po}}(x) = e^{-x},$$

- Gaussian Ensemble: limit ($N \rightarrow \infty$) of prob. space of $N \times N$ matrices $B = (b_{ij})$ with $b_{ij} \in N(0, 1)$ i.i.d. with some measure $P(B) dB$.
 - Gaussian Orthogonal Ensemble (GOE). If B is symmetric. Time-reversal symmetry
 - Gaussian Unitary Ensemble (GUE). If B is unitary. No time-reversal symmetry.
 - Gaussian Symplectic Ensemble (GSE) if B is symplectic. If there is half-integer spin.
- Normalized spacing distribution is

$$\rho_{\text{G}*E}(x) = ce^{-\sigma x^2}$$

The Cat map (briefly)

- Simple model: $A \in SL_2(\mathbb{Z}) : \mathbb{T}^2 \rightarrow \mathbb{T}^2$,

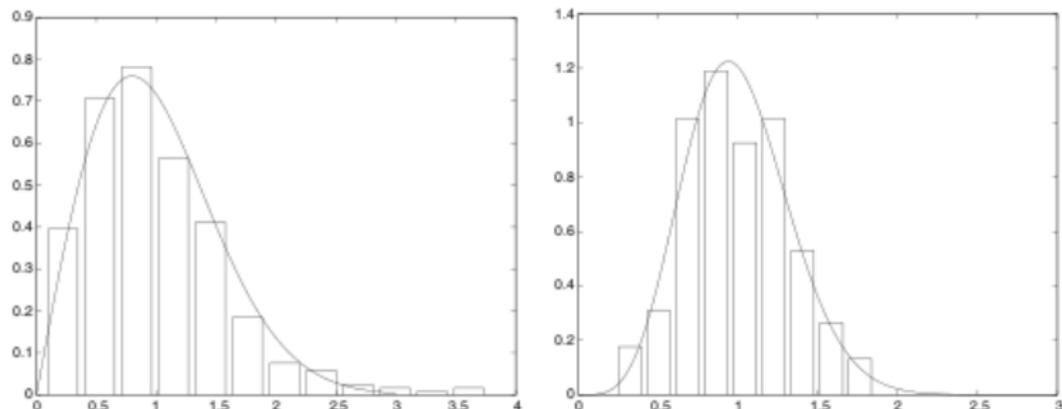
$$\begin{pmatrix} z_1 \\ z_2 \end{pmatrix} \mapsto A \begin{pmatrix} z_1 \\ z_2 \end{pmatrix} = \begin{pmatrix} az_1 + bz_2 \\ cz_1 + dz_2 \end{pmatrix}$$

- Quantization \Rightarrow unitary matrix $U_N(A)$ acting on $L^2(\mathbb{Z}/N\mathbb{Z})$.
- $U_N(A)$ is given by the Weil representation associated to the discriminant form $(\mathbb{Z}/N\mathbb{Z}, x \mapsto Nx^2)$
- Consider action of a group $\Gamma_{A_1, \dots, A_n} = \langle A_1, \dots, A_n \rangle$ with associated $z \in \mathbb{C}[SL_2(\mathbb{Z})]$

$$z = A_1 + A_1^{-1} + A_2 + A_2^{-1} + \dots + A_n + A_n^{-1}$$

- Dynamical properties of the quantum system depends on properties of z .

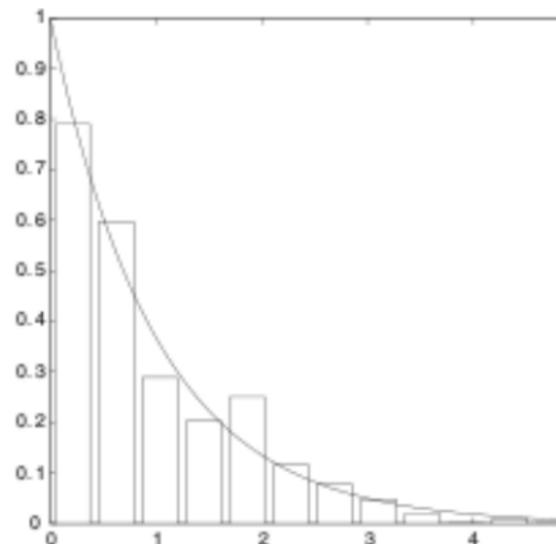
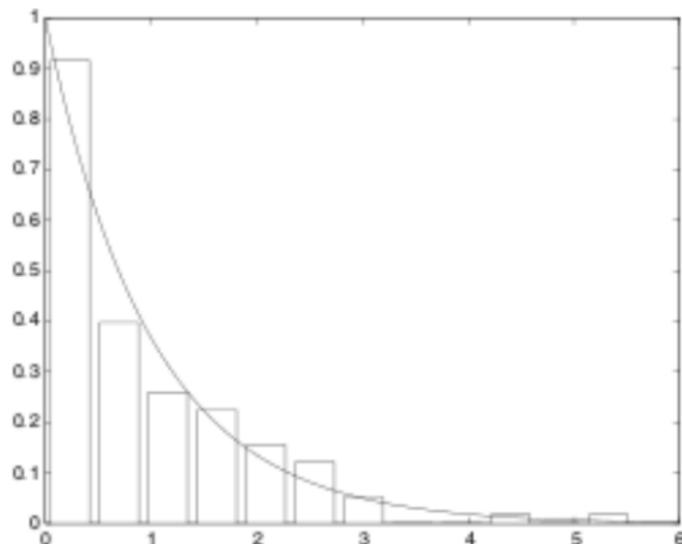
Example: „Random“ Quantum Cat Map



Here z is random (left: GOE, right: GSE).

(Gamburd, Lafferty and Rockmore, „Eigenvalue spacings for quantized cat maps“, 2003)

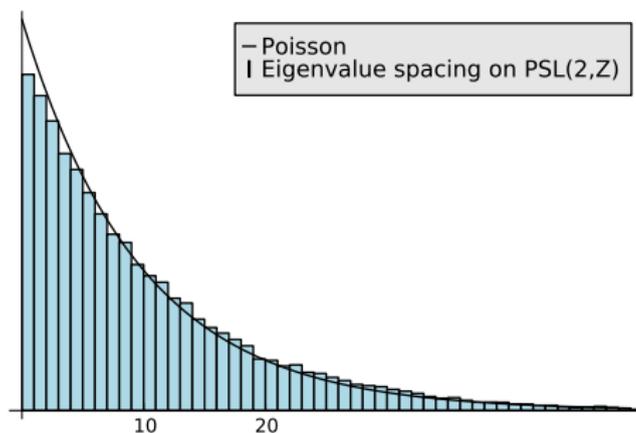
Example: „Arithmetic“ Quantum Cat Map



Here z is a Ramanujan element (Poisson).

(Gamburd, Lafferty and Rockmore, „Eigenvalue spacings for quantized cat maps“, 2003)

Example: The modular group

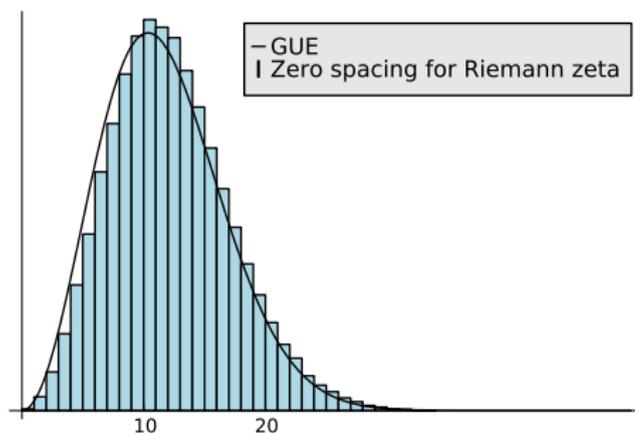


Approx. 50000 eigenvalues of the modular group (data computed by H. Then).

Example: Riemann Zeta

Conjecture

Riemann zeta function zeros are supposed to follow GUE distribution (Montgomery, Dyson)



The first 100000 zeroes of ζ (data computed by Odlyzko)

Maass waveforms in magnetic field

- Set $\bar{\Gamma} = \pi^{-1}(\Gamma) \subset SL(2, \mathbb{R})$ where π is the natural proj.
- The Laplacian with weight/magnetic fieldstrength k (Landau gauge)

$$\Delta_k = y^2 \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} \right) - iky \frac{\partial}{\partial y} = \Delta - iky \frac{\partial}{\partial y}.$$

- Automorphy factor:

$$j_A(z)^k = e^{ik \operatorname{Arg}(cz+d)} = \left(\frac{cz+d}{c\bar{z}+d} \right)^{\frac{k}{2}},$$

$$\sigma_k(A, B) = j_A(Bz)^k j_B(z)^k j_{AB}(z)^{-k} \in \{1, e^{\pm 2\pi i k}\}$$

- Weight k action:

$$\varphi|_k A = j_A(z)^{-k} \varphi(Az), \quad A = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \Gamma$$

Multiplier systems

- Multiplier system $\nu : \Gamma \rightarrow S^1$ satisfies

$$\begin{aligned}\nu(-1) &= e^{-\pi i k} \\ \nu(A)\nu(B) &= \sigma_k(A, B)\nu(AB), \quad \forall A, B \in \bar{\Gamma}.\end{aligned}$$

- Existence of ν is equiv. to existence of $f \in C^\infty(\mathcal{H})$, $f \not\equiv 0$

$$f(Az) = \nu(A)(cz + d)^k \varphi(z), \quad \forall A \in \bar{\Gamma}$$

- Maass waveform $\mathcal{M}(\Gamma, \nu, k) \ni \varphi$:

$$\begin{aligned}(\Delta_k + \lambda)\varphi &= 0, \\ \varphi|_k A &= \nu(A)\varphi, \quad \forall A \in \bar{\Gamma}, \\ \int_{\Gamma \backslash \mathcal{H}} |\varphi|^2 d\mu &< \infty.\end{aligned}$$

Magnetic Maass waveforms

- The multiplier determines phase factors around cusps:

$$v(S_j) = e^{i\alpha_j}$$

(There is no continuous spectrum from cusp nr. j if $\alpha_j \neq 0$)

- If $\Gamma = PSL_2(\mathbb{Z})$ and $k \in \mathbb{R}$ there are 6 possibilities, all related to v_η^{2k} , where

$$v_\eta(A) = \eta(Az) / \eta(z)$$

(independent of $z \in \mathcal{H}$) with Dedekind's eta function:

$$\eta(z) = e\left(\frac{z}{24}\right) \prod_{n \geq 1} (1 - e(nz)), \quad z \in \mathcal{H}.$$

- $v_k(T) = e\left(\frac{1}{12}\right) \Rightarrow \varphi \in \mathcal{M}(\Gamma, k, v_\eta)$ satisfy

$$\varphi(z+1) = e^{\frac{2\pi ik}{12}} \varphi(z)$$

Correspondences on a Surface

- A *correspondence* C of order r on a surface X is a map

$$\begin{aligned} C : X &\rightarrow X^r/S_r \\ z &\mapsto \{z_1, \dots, z_r\} \end{aligned}$$

where $z_j(z)$ are locally isometries (only defined globally as a set)

- If $X = \Gamma \backslash \mathcal{H}$ choose $\delta \in PGL(2, \mathbb{R})$ such that

$$\delta^{-1}\Gamma\delta \cap \Gamma = B$$

has finite index in both Γ and $\delta^{-1}\Gamma\delta$, i.e. δ is in the commensurator of Γ , $\text{Com}(\Gamma)$.

Hecke operators

- Define

$$C_\delta : \Gamma z \rightarrow \{\Gamma \delta \alpha_1 z, \dots, \Gamma \delta \alpha_r z\}$$

where $\Gamma = \cup B \alpha_j$ is a right-coset decomposition.

- This correspondence induces a Hecke operator

$$T_\delta : L^2(X) \rightarrow L^2(X),$$

$$T_\delta f(z) = \sum_{j=1}^r f(\delta \alpha_j z).$$

- T_δ commutes with Δ and generates an algebra which is large if $\text{Com}(\Gamma)$ is large.
- Margulis: $\text{Com}(\Gamma)$ is dense in $PGL(2, \mathbb{R})$ if and only if Γ is arithmetic and otherwise $\text{Com}(\Gamma)/\Gamma$ is a finite group.



Example for $\Gamma = PSL_2(\mathbb{Z}) \subseteq PGL(2, \mathbb{R})$

It can be shown that

$$\begin{aligned} \text{Com}(\Gamma) &\simeq \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in \mathbb{Z}^{2 \times 2}, ad - bc \neq 0 \right\} \\ &= \bigcup \Delta(n) \end{aligned}$$

where

$$\Delta(n) = \{A \in \mathbb{Z}^{2 \times 2} \mid \det A = n\},$$

and $\Delta(n)/\Gamma$ is a finite set.

Classical Hecke Operators

- As generators we can take p prime and $\delta_p = \begin{pmatrix} p & 0 \\ 0 & 1 \end{pmatrix}$ or $\begin{pmatrix} 1 & 0 \\ 0 & p \end{pmatrix}$
- Classical Hecke operator:

$$T_p f(z) = \sum_{j=0}^{p-1} f\left(\frac{z+j}{p}\right) + f(pz)$$

- T_n for n non-prime can be defined by

$$T_m T_n = \sum_{d|(m,n)} d T_{\frac{mn}{d^2}}.$$

- Note that for $\gcd(m, n) = 1$ we get

$$T_m T_n = T_{mn}.$$

Example: T_2

Consider $p = 2$. $\delta = \delta_2 = \begin{pmatrix} 1 & 0 \\ 0 & 2 \end{pmatrix}$ Then

$$B = \Gamma \cap \delta^{-1} \Gamma \delta = \left\{ \begin{pmatrix} a & 2b \\ c & d \end{pmatrix}, ad - bc = 1 \right\} = \Gamma^0(2),$$

$$B \backslash \Gamma = \{1, S, T\} =: \{\alpha_1, \alpha_2, \alpha_3\}.$$

$$R_1 = \delta \alpha_1 = \delta,$$

$$R_2 = \delta \alpha_2 = \delta S$$

$$T_2 f(z) = f\left(\frac{z}{2}\right) + f\left(\frac{z+1}{2}\right) + f(2z).$$

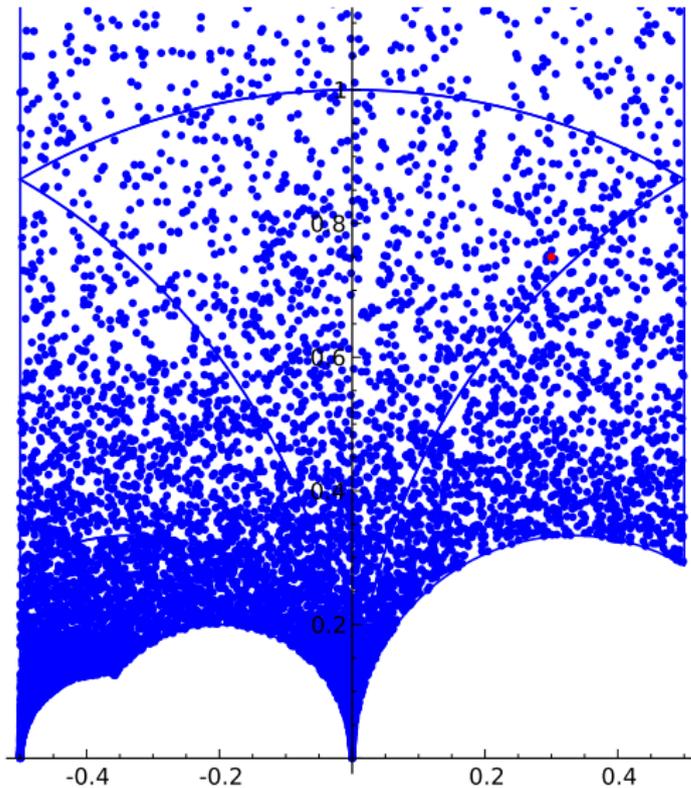
Hecke Points

- Hecke operators are averages over „Hecke points“
- $T_{\rho,z} : f \mapsto \sum f(z_j)$
- We know (Sarnak) that for any $w \in \mathcal{H}$

$$\lim_{\rho \rightarrow \infty} T_{\rho} f(w) = 2\pi \int_{\Gamma \backslash \mathcal{H}} f(z) d\mu(z)$$

(Clozel-Oh-Ullmo, Eskin etc has generalized this to other setting than $SL_2(\mathbb{Z})$)

Hecke Points



$$p = 10007$$



A Commuting Family

- $T_1 = \text{id}$ and $T_{-1}f(z) = f(-\bar{z})$.
- $\mathcal{T} = \{T_n | n \neq 0\}$ is a commuting family of self-adjoint operators on $L^2(\Gamma)$
- They also commute with the Laplacian.
- We can assume that

$$\mathcal{M}(\Gamma) = \bigoplus \mathcal{M}_\lambda(\Gamma)$$

where each $\mathcal{M}_\lambda(\Gamma)$ has a basis of simultaneous eigenfunctions.

Hecke eigenforms

Let $\varphi \in \mathcal{M}(\Gamma)$ be given by

$$\varphi(z) = \sum_{n \neq 0} c(n) \sqrt{y} K_{iR}(2\pi |n| y) e(nx).$$

- $T_{-1}^2 = T_1 \Rightarrow$ eigenvalues 1 and -1 :

$$T_{-1}\varphi(z) = \varepsilon\varphi(z), \quad \varepsilon = \begin{cases} 1, & \varphi \text{ is even,} \\ -1, & \varphi \text{ is odd.} \end{cases}$$

- Note:

$$\begin{aligned} T_{-1}\varphi(z) &= \sum_{n \neq 0} c(n) \sqrt{y} K_{iR}(2\pi |n| y) e(-nx) \\ &= \sum_{n \neq 0} c(-n) \sqrt{y} K_{iR}(2\pi |n| y) e(nx) \\ &= \varepsilon\varphi(z) \Rightarrow c(-n) = \varepsilon c(n) \end{aligned}$$

Hecke Eigenforms

- Define

$$\text{cs}(x) = \begin{cases} \cos(x), & \varepsilon = 1, \\ \sin(x), & \varepsilon = -1. \end{cases}$$

Then

$$\varphi(z) = \sum_{n=1}^{\infty} a(n) \sqrt{y} K_{iR}(2\pi |n| y) \text{cs}(2\pi n x)$$

Hecke Operators and Fourier Coefficients

- We know $a(1) \neq 0 \Rightarrow$ can set $a(1) = 1$.
- Then (T_p replaced by $\frac{1}{\sqrt{p}} T_p$) the p -th Hecke eigenvalue $\lambda_p = a(p)$, i.e.

$$T_p \varphi(z) = a(p) \varphi(z).$$

- Conjectures:
 - Ramanujan-Petersson:

$$|\lambda_p| \leq 2$$

(Holomorphic case: theorem of Deligne)

- Sato-Tate: as $p \rightarrow \infty$ the λ_p are distributed according to semi-circle distribution.

(Holomorphic case: theorems of Conrey-Duke-Farmer, Taylor)

Multiplicity One

Theorem

The common eigenspaces of Δ and T_p are all one-dimensional

Meaning that all problems stemming from high multiplicity (in the spectrum of Δ) vanishes when looking at Hecke eigenforms!

This enables the proof of arithmetic QUE! First key Lemma:

Lemma

If Σ is non-empty subset of geodesics there is a Hecke correspondence C_n and $x_0 \in \Gamma \backslash \mathcal{H} \notin \Sigma$ s.t.

$$C_n x_0 \cap \Sigma = \{x_1\}$$

(For the rest of the proof see the whiteboard!)

What does it mean to compute a Maass form?

- Three algorithms ($\sigma_\varepsilon(\Gamma)$ is an ε -nbhd of the spectrum of Γ):

1 Compute coefficients:

- INPUT: $\varepsilon > 0$ and $R \in \sigma_\varepsilon(\Gamma)$.
- OUTPUT: Sequence of $c(n)$, $1 \leq |n| \leq M$ s.t.

$$\hat{\varphi}(z) = \sum_{|n|=1}^M c(n) K_{iR}(2\pi|n|y) e(nx)$$

satisfies $|\hat{\varphi}(\gamma z) - \varphi(z)| < \varepsilon$ for each $\gamma \in \Gamma$ and $z \in \mathcal{H}$.

2 Test eigenvalue:

- INPUT: $R \in \mathbb{R}$, ε
- OUTPUT: True if $R \in \sigma_\varepsilon(\Gamma)$, False otherwise.

3 Locate eigenvalues:

- INPUT: Interval $I \subset \mathbb{R}$ and $\varepsilon > 0$
- OUTPUT: $\{R_1, R_2, \dots, R_n\}$ s.t. for each $r \in \sigma(\Gamma) \cap I \exists j$ s.t. $|r - R_j| < \varepsilon$.

How do we solve these problems?

There are two types of algorithms (in this area):

- 1 *Rigorous*, meaning that the results are proven to be correct (modulo bugs in the code, quantum effects etc...)
 - Advantage: Can be used to prove theorems.
 - Disadvantages: Slow and hard to implement in many cases.
 - Examples: Verifying eigenvalues using quasi-modes and rigorous estimates (Booker-Strömbergsson-Venkatesh) or locating eigenvalues using the Selberg trace formula (Booker-Strömbergsson).
- 2 *Heuristic*, using independent tests to check results.
 - Advantage: Fast, easy to adapt to many situations.
 - Disadvantage: Could output wrong results.
 - Examples: Algorithms by Hejhal, S., Then, Avelin etc. dealing with different types of functions, groups, multipliers etc.

Overview of the Heuristic Algorithm

- Given a Hecke eigenform $\varphi \in \mathcal{M}(\Gamma; R)$, its Fourier exp.

$$\varphi(z) = \sum_{n=1} c(n) \sqrt{y} K_{iR}(2\pi |n| y) \text{cs}(nx)$$

decays rapidly and we can truncate it: Let

$$\hat{\varphi}(z) = \sum_{1 \leq n \leq M_0} c(n) \sqrt{y} K_{iR}(2\pi |n| y) \text{cs}(nx) = \varphi(z) + [[\varepsilon]].$$

- Treating $\hat{\varphi}(z)$ as a finite Fourier series we use Fourier inversion to obtain the coefficients $c(n)$.
- To get a non-tautological system use automorphic properties of φ , i.e.

$$\varphi(\gamma z) = \varphi(z), \quad \forall \gamma \in \Gamma.$$

- We then arrive at a linear system which we can solve.

Fourier Inversion and Automorphy

Consider a horocycle h at height $0 < Y < Y_0$ and $Q > M_0$.

- Equidistributed pts.: $z_m = x_m + iY$, $1 \leq j \leq Q$ with

$$x_m = \frac{2j-1}{4Q}.$$

- Then

$$c(n) \sqrt{Y} K_{iR}(2\pi |n| Y) = \frac{2}{Q} \sum_{j=1}^Q \varphi(z_m) \text{cs}(-nx_m).$$

- Let $z_m^* \in \Gamma z_m \cap \mathcal{F}_\Gamma$ be the pullback of z_m . Then

$$\begin{aligned} c(n) \sqrt{Y} K_{iR}(2\pi |n| Y) &= \frac{1}{2Q} \sum_{j=1-Q}^Q \varphi(z_m^*) \text{cs}(-nx_m) + [[\varepsilon]] \\ &\simeq \sum_{|n| \leq M_0} V_{nI} c(l) \end{aligned}$$

Algorithm

- Need Y_0 s.t. if $Y < Y_0$ then $\Im z_m^* \geq Y$ for truncation to work with same M_0 ,
- Let Y_0 be the *invariant height* of Γ
- Note: for more than one cusp we use normalizers sending the cusp to infinity.
- For $\Gamma_0(N)$: $Y_0 = \frac{\sqrt{3}}{2N}$.

Tests

- Given a pair R, Y we get a vector $C(Y, R) = (c(1), \dots, c(M_0))$
 - Real analytic in R away from zeros of $K_{iR}(2\pi|n|Y)$.
 - Not necessarily continuous in Y
- Question: How do we know if R is an eigenvalue?
 - $C(Y, R)$ should be independent of $Y < Y_0$.
 - $c(m)c(n) = c(mn)$ for $(m, n) = 1$.
- These kind of tests are also used to locate eigenvalues.

Illustration of location

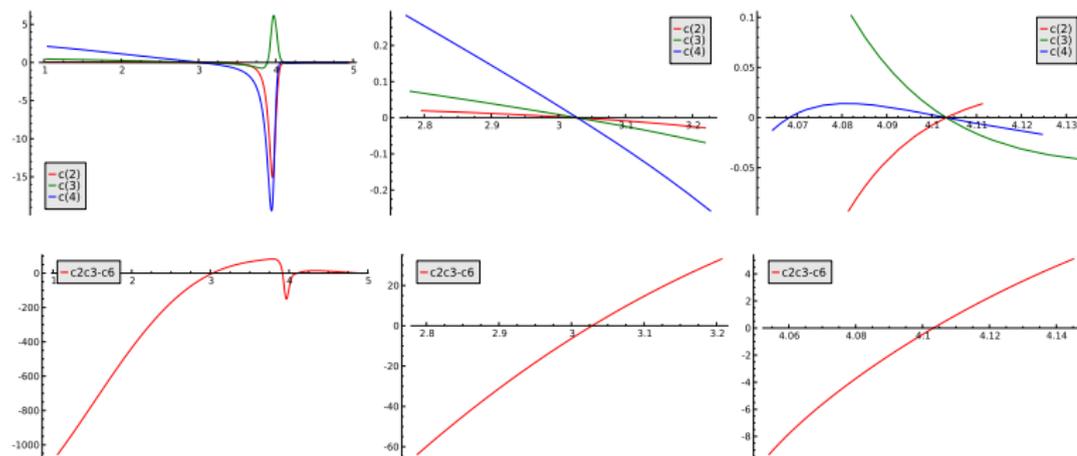
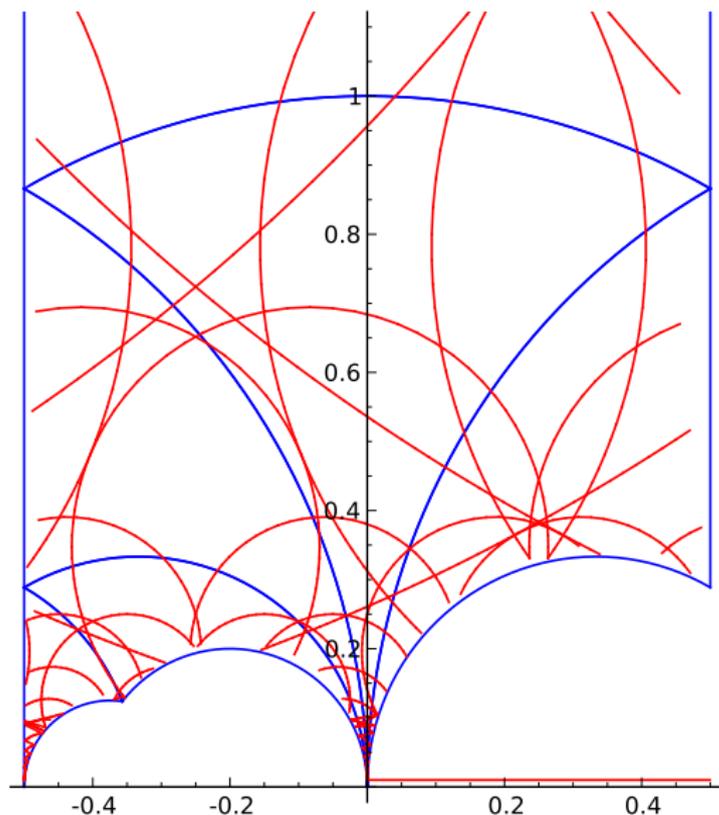


Illustration of the pullback



$$Y = 0.01$$

$$I = 50.00$$

Algorithm: The linear system

- Subtracting diagonal term results in a linear (stable) $M_0 \times M_0$ system

$$\sum_{1 \leq n \leq M_0} \tilde{V}_{nl} c(l) = 0, \quad 1 \leq |n| \leq M_0 \Leftrightarrow$$

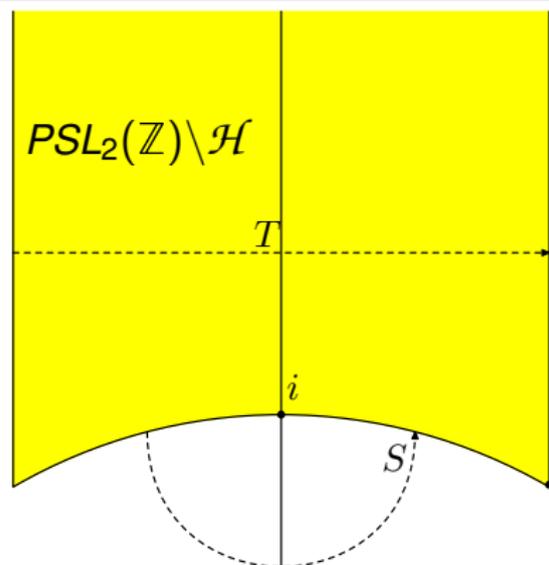
$$\tilde{V}C = 0$$

- Since φ is a Hecke eigenform we can set $c(1) = 1$ and delete the corresponding equation, i.e. we get

$$\tilde{V}'C = W$$

where \tilde{V}' is $M_0 - 1 \times M_0 - 1$ lower right part of \tilde{V} and $W = (-V_{n1})_{1 \leq n \leq M_0}$.

Pullback for the modular group:



$$S = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}, z \mapsto -\frac{1}{z}$$

$$T = \begin{pmatrix} 1 & 1 \\ 0 & 1 \end{pmatrix}, z \mapsto z + 1$$

ALGORITHM: INPUT:

$$z_0 = x_0 + iy, A = id$$

- 1 At step j : $z = x + iy$
- 2 If $|x| > \frac{1}{2}$ let $n \in \mathbb{Z}$ s.t.
 $|x - n| < \frac{1}{2}$:
 - $z \mapsto z - n$
 - $A \mapsto TA$
- 3 If $|z| < 1$:
 - $z \mapsto -\frac{1}{z}$
 - $A \mapsto SA$
 - Continue at (2)
- 4 OUTPUT: z, A .

Pullback for subgroups

- Let $\Gamma \subseteq PSL_2(\mathbb{Z})$ of index n , i.e.

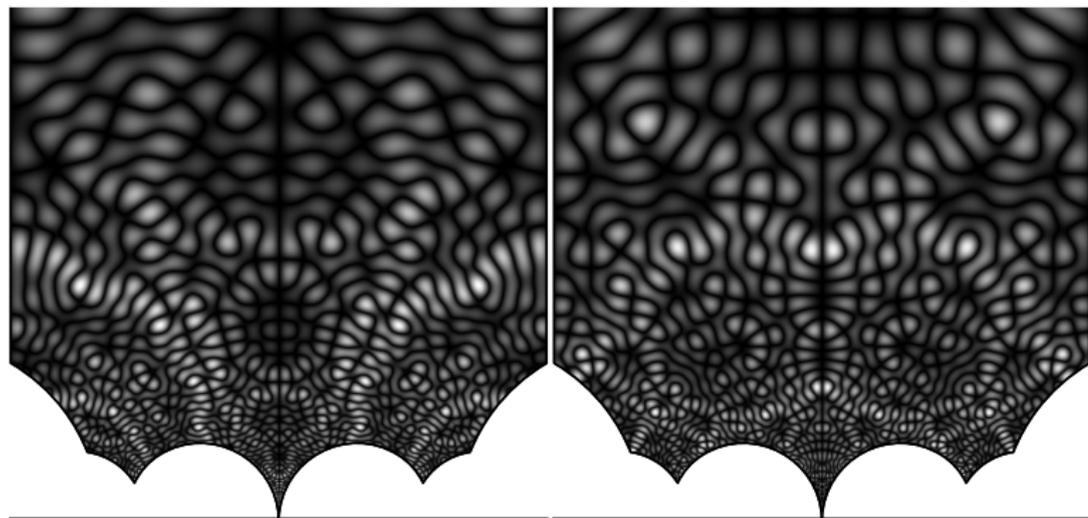
$$\begin{aligned}\Gamma \backslash PSL_2(\mathbb{Z}) &= \{V_1, \dots, V_n\}, \quad \text{or} \\ PSL_2(\mathbb{Z}) &= \Gamma V_1 \sqcup \Gamma V_2 \sqcup \dots \sqcup \Gamma V_n.\end{aligned}$$

- A fundamental domain for Γ is

$$\mathcal{F}_\Gamma = \cup V_j \mathcal{F}_{PSL_2(\mathbb{Z})}$$

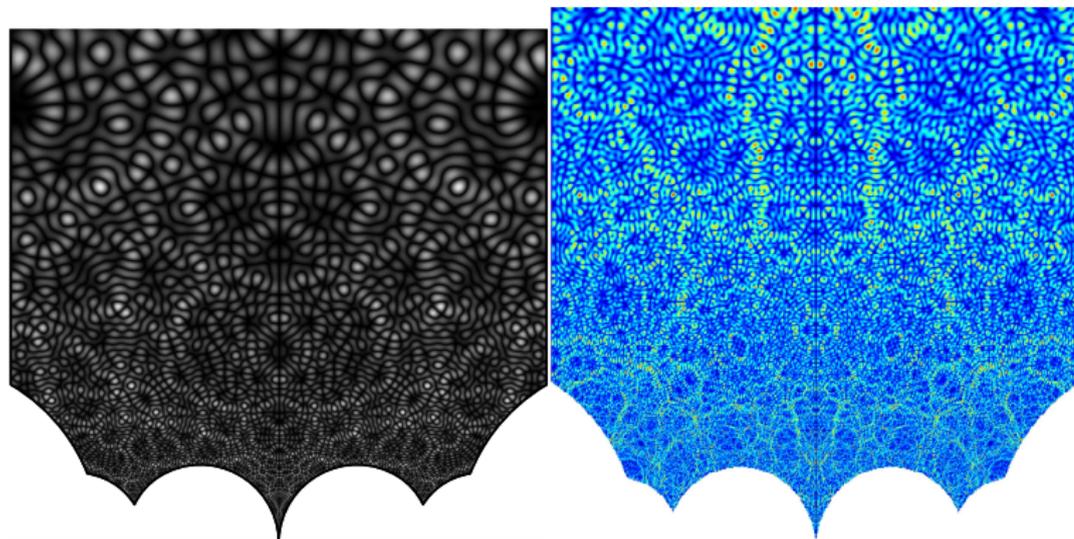
- Let $z \in \mathcal{H}$.
- First get pullback to $PSL_2(\mathbb{Z})$: $\tilde{z} = Az$.
- If $A^{-1} \in \Gamma V_k$ then $A^{-1} V_k^{-1} \in \Gamma \Rightarrow V_k A \in \Gamma$ and $z^* = V_k A z \in \mathcal{F}_\Gamma$.

Maassforms for $\Gamma_0(7)$



$R \simeq 50 \Rightarrow \lambda \simeq 2500$ (approx. 1500th e.v.)

Examples of Maass forms for $\Gamma_0(7)$



Left: $R \simeq 100$, $\lambda \simeq 10000$, approx. 6400th e.v.

Right: $R \simeq 200$, $\lambda \simeq 40000$, approx. 26000th e.v.

Holomorphic Modular forms

Automorphy factor and slash-action of weight k .

$$J_A(z)^k = (cz + d)^k,$$
$$f|_k A(z) = J_A(z)^{-k} f(Az), \quad A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}.$$

Note: $A'(z) = J_A(z)^{-2}$.

Definition (A holomorphic) modular form is)

a holomorphic function $f : \mathcal{H} \rightarrow \mathbb{C}$ satisfying

- 1 $f|_k A = f(z), \forall A \in \Gamma$
- 2 $f(z)$ is holomorphic at the cusps of Γ

Note: If $-1_2 \in \Gamma$ then

$$f|_k -1_2 = (-1)^{-k} f(z) = f(z)$$

hence $k \in 2\mathbb{Z}$.

Spaces of Modular forms

- $\mathcal{M}_k(\Gamma)$ - the space of modular forms of weight k
- $\mathcal{S}_k(\Gamma)$ - the subspace of cusp forms.
- Dimension \mathcal{M}_k and $\mathcal{S}_k = O(k)$.
- The algebra $\mathcal{M}_*(\Gamma) = \bigoplus \mathcal{M}_k(\Gamma)$ is finitely generated.
- Alternative interpretations of \mathcal{M}_{2k} :
 - Holomorphic k -fold differentials on $\Gamma \backslash \mathcal{H}$ via
 $f \in \mathcal{M}_{2k}(\Gamma) \mapsto dZ^k = f(z)(dz)^k$
 - Holomorphic sections of the k -th tensor power of the canonical line bundle on $\Gamma \backslash \mathcal{H}$
 (Note: Line bundles on $\Gamma \backslash \mathcal{H}$ are in 1-1 corr. with the Automorphy factors:

$$(A, z) \mapsto J_A(z)^k.$$

- Analogues of large eigenvalues are large weights.
- $f \in \mathcal{S}_k(\Gamma) \Rightarrow F = y^{\frac{k}{2}} f \in \mathcal{M}(\Gamma; \lambda_k, k)$, $\lambda_k = \frac{k}{2} \left(\frac{k}{2} - 1 \right)$.

Holomorphic QUE

For $f \in \mathcal{S}_k(\Gamma)$ define

$$v_f := |f(z)|^2 y^k d\mu(z)$$

- Since $\dim \mathcal{S}_k(\Gamma) \sim ck$ the analogue of general QUE fails.
- However, an arithmetic analogue holds:

Theorem (Holowinsky-Sound)

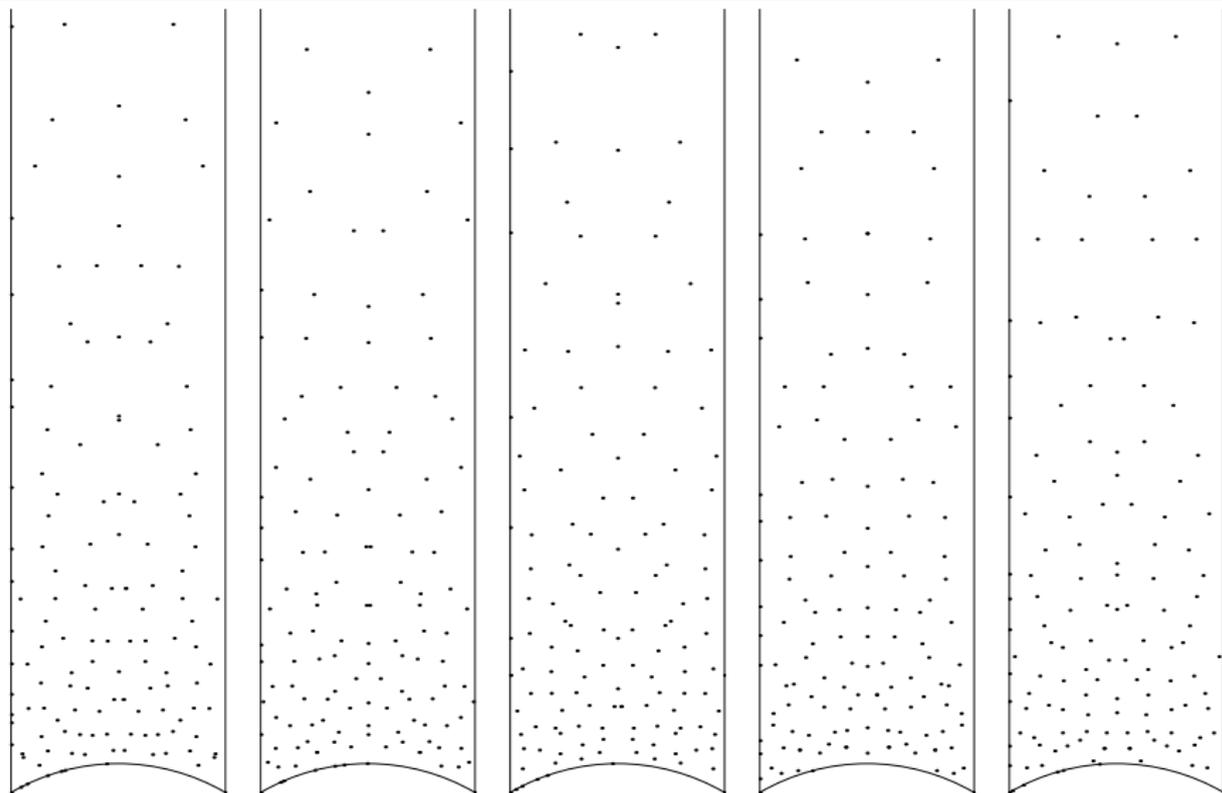
If Γ is non-co-compact and $f \in \mathcal{S}_k(\Gamma)$ is a Hecke eigenform of weight k , then as $k \rightarrow \infty$:

- 1 $v_f \rightarrow \frac{1}{\mu(\Gamma \backslash \mathcal{H})} d\mu$
- 2 The zeros of f becomes equidistributed in $\Gamma \backslash \mathcal{H}$ with respect to $\frac{1}{\mu(\Gamma \backslash \mathcal{H})} d\mu$.

Remarks

- Soundararajan's method involves weak subconvexity for degree 8 L-functions of automorphic representations on $GL(n)$.
- His method applies to both compact and noncompact surfaces but misses $O_\varepsilon(k^\varepsilon)$ forms.
- Holowinsky's method uses the Fourier expansions and therefore applies only to non-compact surfaces.
- He also misses $O_\varepsilon(k^\varepsilon)$ forms.
- However, the sets of functions which are missed by these two methods are disjoint!

Illustration of zeros: $k = 2000$, $\dim S_k = 166$



Remarks on zeros of $f \in \mathcal{S}_k(PSL_2(\mathbb{Z}))$

$$f = \sum_{n=1}^{\infty} a(n) e(nz), \quad a(n) \in \mathbb{R}$$

- f has $\frac{k}{12} + O(1)$ zeros.
- then the zeros of f are symmetric w.r.t. $i\mathbb{R}$, i.e.
 $f(z) = 0 \Leftrightarrow f(-\bar{z}) = 0$.
- f is real-valued on $\delta_1 : \Re z = 0$ and $\delta_2 : \Re z = \frac{1}{2}$.
- $z^{\frac{k}{2}} f$ is real-valued on $\delta_3 : \{|z| = 1, 0 \leq x \leq \frac{1}{2}\}$.
- By Sarnak-Gosh: The number of zeros of f on $\delta_1 \cup \delta_2 \cup \delta_3$ is $\gg_{\varepsilon} k^{\left(\frac{1}{4} - \frac{1}{60} - \varepsilon\right)}$ for any $\varepsilon > 0$ (Conj: $\sqrt{k} \ln k$).
- For $Y \gg \sqrt{k \ln k}$ almost all zeros of height $\geq Y$ are on $\delta_1 \cup \delta_2$.

Some Further Reading

- Bogomolny, Georgeot , Giannoni and Schmit, „Quantum Chaos on Constant Negative Surfaces“, Chaos, Solitons and Fractals, 1995.
- Sarnak, Arithmetic Quantum Chaos, 1993 (Lecture notes).
- Sarnak, Recent Progress on QUE, BAMS, 2011.
- Zelditch, Recent Developments in Mathematical Quantum Chaos, 2009 (Lecture notes).
- Menezes, Jar e Silva and de Aguiar „Numerical experiments on quantum chaotic billiards“, AIP Chaos 2007.
- Soundarajan, „Quantum unique ergodicity for $PSL_2(\mathbb{Z})$ “, Ann. Math. 172, 2010.
- Holowinsky and Soundararajan, „Maass equidistribution of Hecke eigenfunctions“, Ann. Math. 172, 2010.